EXPERIMENTAL AND NUMERICAL STUDY OF PHASE TRANSITION OF LiFePO₄ MATERIAL IN LITHIUM ION BATTERIES

by

MD NOOR E ALAM SIDDIQUE

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Abstract

EXPERIMENTAL AND NUMERICAL STUDY OF PHASE TRANSITION OF LIFePO₄ MATERIAL IN LITHIUM ION BATTERIES

Md Noor E Alam Siddique, PhD

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Supervising Professor: Fugiang Liu

Phase transition behavior of LiFePO₄ material has been studied in this work. During electrochemical charge/discharge processes, LiFePO₄ transforms into FePO₄ and this electrochemically driven phase transition of the two-phase system results in a potential plateau in a battery discharge curve. Besides, battery performance, especially under high rates, depends critically on this two-phase transition. However, this phase transition mechanism in the LiFePO₄ crystal structure has yet not been understood in details. Developing better understanding is essential for designing high performing, safe and stable batteries.

Currently available phase transition models for LiFePO₄, such as the classical 'core-shrinking model' and recently the 'domino cascade model', have shed lights on the phase transition mechanism, however, with somewhat contradicting conclusions. The former predicted a continuous phase transition between LiFePO₄ and FePO₄ within an individual particle. In contrast, the latter, based on the fact of anisotropic diffusion of Li, provided a microscopically heterogeneous picture of phase transition: phase transition is abrupt, and de-lithiated and lithiated particles co-exist in a battery electrode.

To achieve a better understanding of the two-phase transition phenomena, two specially designed sophisticated methods, i.e., in-situ Raman spectroscopy and in-situ Xray Diffraction (XRD), have been developed in this work. The particle level phase transition of LiFePO₄ was probed in the in-situ Raman study. Under slow-rate discharge/charge, it was found that the particles were either fully transformed or untransformed indicating that the phase transition in LiFePO₄ material was not uniform at the particle level. Electronic conductivity and local electrode microstructure determined the preferred sequence of phase transition of the particles. This study also revealed that non-equilibrium phase transition occurred as a result of delayed phase transition. The insitu XRD study attempted to investigate the bulk level phase transition behavior of LiFePO₄ and was conducted at four different discharge rates. Analysis of the highresolution diffraction patterns showed a gradual and continuous change in phase transition. Again, a delayed phase transition was observed at high-rate discharge due to the non-equilibrium pathway. From both in-situ Raman and in-situ XRD experiments, only two stoichiometric compounds (LiFePO₄ and FePO₄) was detected, indicating that during electrochemical charge/discharge processes once Li ion de-intercalates/intercalates, it immediately forms the end-members and no intermediate compounds exist.

Particle size and morphology of LiFePO₄ were also studied using numerical models in the phase transition behavior. At a slow discharge rate, plate shaped particles showed higher capacity than spherical particles and the opposite trend was observed at high discharge rates above some critical discharge rates. Big particles and high discharge rates induced non-equilibrium phase transition in the LiFePO₄ particles which expands the miscibility gap between the end-members and gave sharp narrow two-phase interfaces.

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Chapter 1

Introduction

Extensive use of energy storage or batteries started with revolution in development of portable devices like cell phones, cameras, personal notebooks etc. Some of these portable devices can use primary batteries or secondary batteries that are also known as rechargeable batteries. One of the widely known rechargeable batteries is the led acid batteries. Later on, small sized nickel-cadmium, nickel-metal hydride batteries were also developed which occupy a huge consumer market. With the development of these secondary batteries, and also different consumer electronic devices, the need for batteries with high energy density becomes more and more prominent to meet the continuous demand for smaller, light weight, reliable, and high performing devices. The first breakthrough to meet these demands came with the development of lithium ion batteries (LIB) in the 90's. Since then, the demand for LIBs increased exponentially and by now already crossed \$13000 million value [1]. The present and future demand scenario is compared in Fig. 1.1, where one most potential application is either plug-in hybrid electric vehicles (PHEV, HEV) or electric vehicles (EV).

LIBs first came with the promise of having high energy density. The energy density of different secondary rechargeable batteries is compared in Fig. 1.2. As is shown, LIBs show higher volumetric and gravimetric energy density than most of the other rechargeable batteries. Higher volumetric and gravimetric energy density enables smaller and lighter batteries. So, this unique battery system becomes an automatic potential choice for EV technologies. For this reason, scientists and engineers focused on this battery system within the last two decades and published thousands of research papers. The research areas include development of new electrode material, electrolyte

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solution, and separator material, optimizing the battery parameters, battery design, safety and durability and obviously reducing cost.

Figure 1.1: Electric vehicle sales and lithium demand for electric vehicle batteries from 2008 to 2020 [2].

The battery technology can be advanced through improving material, design and employing better battery management practices. But without improving the material itself the energy storage technology cannot move forward. This LIB technology is one such example where new material system gave high performing, reliable energy storage system with high energy density. To further improve the charge/discharge capacities and reliability of LIBs, the materials used have to be advanced. Discovering new electrode materials and understanding their structure-performance relationship are keys to the above-mentioned improvement.





When LIBs was first introduced, the cathode material used was mostly transition metal oxide materials. And then one breakthrough in cathode materials came with the introduction of LiFePO₄ of olivine structure [4]. This material is safe, stable, cheap, easy to synthesize and also able to deliver high capacity. The poor electronic conductivity of this material was a problem initially and later on it was fixed using different techniques and approaches.

LiFePO₄ gives a very flat potential plateau during discharge due to the presence of a two phase system. But at high discharge rates, this material shows reduced capacity and reduced potential due to partly lithium ion diffusion problems through anisotropic crystallographic channel and partly due to the unique phase transition phenomena. Slow discharge leads to the equilibrium phase transition (between lithium rich LiFePO₄ and lithium poor FePO₄) where Li diffusion through anisotropic crystallographic channel does not necessarily compete with the Li flux that comes from the electrochemical discharge rates. That results in a stable potential and good capacity during the equilibrium phase transition. On the other hand, non-equilibrium phase transition could appear at higher discharge rates, which is complicate because of the anisotropic Li diffusion, Li ion migration difficulty through porous network and the presence of a narrow interface between end members (LiFePO₄ and FePO₄) with very limited Li ion diffusivity.

Apart from all these issues, the phase transition phenomena in LiFePO₄ material, both equilibrium and non-equilibrium are not well understood so far. There exist many debates on several key issues regarding this phase transition mechanism. This very intrigue and important phase transition phenomenon is the motivation to study further the phase transition mechanism, which forms the central theme of this study. This work thus strives to clarify the phase transition mechanism and also to understand the particle level phenomena of LiFePO₄ materials.

This dissertation begins with a very general description of LIB system together with introduction on battery materials and their selection criteria in Chapter 2. Chapter 2 also includes description of LiFePO₄ material, its phase transition models and representative techniques to study the phase transition, and finally concludes with a list of objectives of this study. Chapter 3 describes the LiFePO₄ synthesis method and the characterization tools and most importantly details of electrode and battery making procedures.

The two main experimental techniques, i.e., in-situ Raman spectroscopy and insitu XRD, to study the phase transition phenomena in LIBs containing LiFePO₄ are discussed in details in Chapter 4 and Chapter 5. Due to experimental limitation every aspect of the phase transition behavior cannot be studied experimentally and hence numerical study on individual LiFePO₄ particles is included in Chapter 6. And lastly Chapter 7 summarizes the key findings of this work and proposes future work.

4

Chapter 2

Background Information and Objective of the Study

2.1 Introduction

LIBs have some basic components: anode, cathode, electrolyte, and separator. Extensive amount of research has already been devoted to improvement of performance of LIBs through developing electrode materials [5-9], reducing particle size [10] and optimizing shape [7, 11, 12] of active material particles of the electrode, and using different additives [13-15].

LiFePO₄, one prominent cathode material that shows a stable plateau of potential during discharge due to the unique phase transition through movement of the two-phase interface [16]. In this chapter, details of the LIB system, materials, their attributes and particularly details of LiFePO₄ materials, its phase transition behavior are discussed. Finally, based on existing study on this material and theoretical knowledge, a scope of research will be identified.

2.2 Li ion Battery System

2.2.1 Battery components

In the schematic illustration of Fig. 2.1, during charge lithium ions come out of the cathode (LiFePO₄), and travel through the electrolyte and separator (which should be permeable to lithium ion), and become intercalated into the anode. During discharge, the above processes reverse. In this particular case of Fig. 2.1, the anode is made of graphitic carbon coated on top of a copper foil along with some binder (PVDF) and the cathode is LiFePO₄-based material coated on an aluminum foil along with conducting additives such as carbon.

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Figure 2.1: Schematic illustration of Li ion battery system including Li ion movement

during charging and discharging.

The battery reactions can be expressed as follows:

Anode reaction: $6C + xLi^+ + xe^- \leftrightarrow C_6Li_x$

Cathode reaction: LiFePO₄ \leftrightarrow Li_(1-x)FePO₄ + xLi⁺ + xe⁻

where the forward direction (from left to right) indicates charge reaction and the reverse direction indicates discharge reaction.

2.2.2 Performance Parameters and Selection Criteria of Battery Materials

Performance of the energy storage systems is relying on several parameters and also depending on the applications, though some specific parameters become more important comparing to others. These parameters include the potential, power the battery delivers to the device, storage capacity, columbic efficiency on energy discharge, charge/discharge rate, etc. The selection of battery materials significantly depends on open circuit voltage (OCV) of the cell. The OCV depends directly on the chemical potential of the electrode materials and is described as [17]

$$V_{OC} = \frac{\left(\mu_A - \mu_C\right)}{F}$$
....2.1

where μ_A and μ_C are the chemical potentials of the anode and cathode materials respectively and *F* is the Faraday constant. In a battery electrode, due to low conductivity of the solid phases and liquid electrolyte phase, as well as slow electrochemical reactions, the voltage V_{dis} obtained during discharge can be decreased by an amount of η_{dis} which is called overpotential [17].

$$\eta_{dis} = V_{OC} - V_{dis} \dots 2.2$$

Similarly, during a charge process a certain overpotential η_{ch} has to be applied together with the OCV and hence the overpotential can be expressed as

where V_{ch} is the potential applied when charging the cell.

In a complete discharge process, the effective energy density stored in the cell is QV/W, where Q is the amount of charge transferred and W is the weight of the cell. During a galvanostatic discharge, near the end of discharge process diffusion of Li ions through electrode materials and Li ion migration through electrolyte limits the capacity (Q/W) which will appear lower than the theoretical capacity.

Battery performance is characterized with two kind of efficiency such as energy efficiency and columbic efficiency. The energy efficiency of a lithium ion cell is [17]

EnergyEfficiency =
$$\begin{pmatrix} \int_{0}^{Q} V_{dis} dq \\ \int_{0}^{Q} V_{ch} dq \end{pmatrix}$$
.....2.4

And the columbic efficiency can be expressed as

$$Columbic Efficiency = \begin{pmatrix} Q_{dis} \\ Q_{ch} \end{pmatrix} \dots 2.5$$

Several cells are connected in series to get a desired voltage *V* and in parallel to get the desired capacity *Q* for specific application.

The open circuit potential can be chosen according to Eq. 2.1 by changing the electrode material. But the electrode material cannot be randomly chosen as the electrolyte might decompose in the operating potentials if the potential is above or below the stability potential window of the electrolyte. The potential window of the electrolyte system is

$$E_g = E_{LUMO} - E_{HOMO} \dots 2.6$$

where the LUMO and HOMO are respectively the lowest unoccupied molecular orbital and highest occupied molecular orbital of the electrolyte. Fermi energy level of the anode E_{FA} has to be lower than LUMO and the Fermi energy level of the cathode E_{FC} has to be higher than the HOMO to fulfill the thermodynamic stability criteria. If the E_{FC} lies above the electrolyte LUMO, the electrolyte will be reduced and if the E_{FA} lies below the electrolyte HOMO, the electrolyte will be oxidized which results in the following thermodynamic stability criterion.

$$eV_{OC} = E_{FA} - E_{FC} \le E_g$$
.....2.7



Figure 2.2: (a) Thermodynamic stability criteria of the electrolyte, (b) Selection criteria for electrode materials based on the thermodynamic stability criteria of the electrolyte [18].

The thermodynamic stability criterion for electrolyte is shown in Fig. 2.2(a) and Fermi energy level for different electrode materials in a typical non-aqueous electrolyte system is shown in Fig. 2.2(b).

2.2.3 Battery Materials

2.2.3.1 Anode and Electrolyte

For the anode, mostly carbon based materials are used, but some other materials like lithium foil, Si [19], Si-C composite [20-22], SnO₂ [23], Li₄Ti₅O₁₂ [24-26] are also used for higher capacity, stability and better performance. Electrolyte systems for LIBs include lithium salts like LiPF₆ [27], LiAsF₆ [28], LiClO₄ [28], LiBF₄ [28], LiCF₃SO₃ [29] dissolved in organic solvents such as ethylene carbonate (EC) [28, 30], dimethyl carbonate (DMC) [30, 31], diethyl carbonate (DEC) [28], etc.

2.2.3.2 Cathode

Cathode materials include oxides, phosphates, sulpho-phosphates that can undergo oxidation and reduction during charge and discharge reactions. A very detailed discussion on cathode materials is published by M Stanley Whittingham [32]. Lithium transition metal oxide such as $LiCoO_2$ [5, 33-38], $LiNiO_2$ [9, 39, 40], $LiMn_2O_4$ [26, 30, 41-44], LiV_3O_8 [45], $Li(Ni,Mn,Co)O_2$ [8, 46], $Li(Ni_{0.8}Co_{0.15}Al_{0.05})O_2$ [6] etc, is a type of material that is very common in commercial LIBs. $LiCoO_2$ is one of the successful commercial cathode materials, but still it has some issues that make it less attractive. $LiCoO_2$ is not very stable at overcharged condition as cobalt dissolution in the electrolyte is believed to happen at a high potential [34]. Also a sharp change in the lattice parameters with lithium intercalation [35] causes stress and micro-cracking of the electrode materials [47]. In addition, cobalt in $LiCoO_2$ is expensive compared to other transition metals such as nickel, manganese, iron, etc. Another transition metal oxide $LiNiO_2$ has higher energy density [48] than $LiCoO_2$, but is less stable [39] and less ordered material [40]. $LiMn_2O_4$ [26] has also been attempted as cathode material which is safer and cheaper than $LiCoO_2$, but shows less capacity. V_2O_5 [49] and LiV_3O_8 [50] are also used as cathode materials. They show higher capacity but lower voltage (3 V or less).

Another class of cathode materials is $LiMPO_4$ [51] with olivine structure which includes $LiFePO_4$, $LiMnPO_4$ [52], $LiCoPO_4$ [53], $LiMn_{0.25}Fe_{0.75}PO_4$ [54, 55], $LiFe_{1-x}Co_xPO_4$ [56], and $Li_3V_2(PO_4)_3$ [57] etc. In this class, $LiFePO_4$ is the most promising one.

Some cathode materials might decompose at overcharge condition or at elevated operating temperatures and these reactions are exothermic which eventually cause the battery temperature to rise. This higher temperature promotes the reactions of cathode materials with non-aqueous solvents (EC, DMC, PC etc) which generate more heat and the temperature continue to rise and eventually at some point the battery might fail and explode and fire hazard can be created. This situation is called the thermal runaway. The thermal property of some cathode materials is compared in Fig. 2.3 [58] and from this figure, it is observed that LiFePO₄ and LiMn₂O₄ are two most safe cathode materials. Besides, LiFePO₄ is cheap and most safe material and this is one good reason that people are motivated to study LiFePO₄.



EC:PC:DMC\LiPF6

Figure 2.3: Comparison of thermal runaway properties of different cathode materials for

LIBs [58].

2.3 LiFePO₄ Materials

Olivine type LiFePO₄ was first described as the battery electrode material by Professor Goodenough [4]. It is relatively cheap, has very stable potential plateau (3.5 V vs Li⁺/Li theoretically) during discharge [59-61], and theoretical capacity of 170 mAh/g [4, 62] and good stability against thermal runaway [63]. Moreover it is observed in Fig. 2.2(b) that LiFePO₄ material is stable in the electrolyte potential window according to the thermodynamic stability criteria. Apart from having all attractive properties, its intrinsic electronic conductivity is very poor which is of the order of 10^{-9} S/cm² [32]. Several strategies have been adopted to increase the conductivity of LiFePO₄ such as conductive carbon coating [64-68], particle size reduction [48], metal powder addition [69, 70], doping with supervalent ions [65, 67, 71, 72], etc.

2.3.1 Structure of LiFePO₄



Figure 2.4: a) Crystal structure of LiFePO₄, b) a layer of LiFePO₄ which consists of FeO_6 octahedra and PO₄ tetrahedra [73].

LiFePO₄ is an olivine structure (*Pnma*) with an orthorhombic unit cell. The crystal structure can be better explained with the help of the illustration in Fig. 2.4. LiFePO₄ structure has some parallel sheets made of FeO₆ octahedra which have corner sharing

connections to each other. These $[FeO_4]_n$ sheets are parallel to 'bc' plane which connect to other $[FeO_4]_n$ sheet with the help of PO₄ tetrahedron along 'a' direction making a 3D structure which has lithium tunnel parallel to the 'b' direction. PO₄ tetrahedron has a common edge with FeO₆ octahedron. Fig. 2.4(a) shows the 3D structure of LiFePO₄ and the individual $[FeO_4]_n$ sheet is shown in Fig. 2.4(b). Lattice parameters of LiFePO₄ (discharged state) are *a*=10.3375 Å, *b*=6.0112 Å, and *c*=4.6950 Å [74] and the lattice parameters for FePO₄ (charged state) are *a*=9.7599 Å, *b*=5.7519 Å, and *c*=4.7560 Å [74]. 2.3.2 Thermodynamics of LiFePO₄

Thermodynamics of a two-phase system is introduced here to illustrate the twophase transformation phenomenon. Gibbs free energy versus composition plot is shown in Fig. 2.5. At room temperature the free energy curves of both α (FePO₄) and β (LiFePO₄) have minima at some particular compositions as shown in Fig. 2.5. In the composition range from α_e to β_e , the criterion that requires minimum Gibbs free energy of the system leads to a two-phase mixture typically at room temperature.



Figure 2.5: Schematic representation of Gibbs free energy with the Li ion content and associated two-phase mixture region.

During discharge of a LIB, Li ions diffuse into the FePO₄ particles and after reaching the composition of X₁ in Fig. 2.6 the two phase mixture appears and it continues until the composition X₂ is achieved. Chemical potentials are obtained from the common tangent of the Gibbs energy curves (Fig. 2.5). Equilibrium potential of the battery can be calculated according to the Eq. 2.1 where $\mu_{Li(c)}$ and $\mu_{Li(a)}$ represent the chemical potentials of the cathode and anode materials respectively. Anode material is usually a single phase material throughout the Li ion intercalation and hence the chemical potential is usually constant. Chemical potential of the cathode material is plotted in the upper part of Fig.2.6 and thermodynamic battery potential (OCV) derived from the Gibbs free energy curve is placed in the bottom part of the Fig. 2.6 which was described by Eq. 2.1.



Figure 2.6: Schematic representation of Gibbs free energy of the α - β two-phase

system and associated potential profile during discharge.

2.3.3 Phase Transition Models

During discharge, Li ion diffuses into FePO₄ (Heterosite) particles which transforms into a new phase LiFePO₄ (Triphylite) and during the charge process the reverse phase transition takes place. This phase transition relies on diffusion of Li ions through the crystal structure of active materials. Diffusion of Li ions through the LiFePO₄ structure is not isotropic but anisotropic along the preferred 'b' direction [73]. This makes Li ion diffusivity higher along the 'b' direction [73] which is due to the 1D tunnel shown in the structure of LiFePO₄ in Fig. 2.4. This anisotropic Li diffusion makes the phase transformation more complicated. Several phase transition models such as coreshrinking model [4, 75], domino-cascade model [73], mosaic model [76, 77], mushy zone approach [7], single phase kinetic model [78] have been postulated by different research groups to describe phase transition phenomena of LiFePO4 materials. And in many aspects these models somehow contradicts with each other which will be highlighted in next few sections. The two-phase models can be broadly categorized into two groups such as the core-shrinking model and domino-cascade model which will be discussed in detail in the following sections.

2.3.3.1 Core Shrinking Model

One of the pioneering works regarding phase transformation study of LiFePO₄ particles was done by Venkat Srinivasan and John Newman [75] which is known as the so-called core-shrinking model or core-shell model. This model was first proposed by Padhi et. al. [4] and later was advanced by Srinivasan and Newman [75], and L Laffont and C Delacourt et. al. [79].

Venkat Srinivasan and John Newman considered phase transformation in spherical particles. Their approach is schematically illustrated in Fig. 2.7. During discharge, Li ions diffuse into the particles starting from the surface which initially creates

 Li_yFePO_4 where y is a very small value. Then a thin shell of $Li_{1-x}FePO_4$ is formed where x is close to zero which is a distinctively different phase than the initial FePO_4 phase. With continuing discharge (Li intercalation), the shell continues to grow until the whole particle becomes $Li_{1-x}FePO_4$ and later on completely transforms into $LiFePO_4$ at the completion of discharge. So, initial single phase transforms through a two-phase system with diffusion of Li ions through a thin moving interface. During the charge, the Li diffusion direction reverses causing a Li-poor shell and Li-rich core where the shell grows inward to finally form FePO_4.



Figure 2.7: Illustration of the shrinking-core model during discharge. [75]

The key attributes of this core-shell model or core-shrinking model are as follows:

 The phase transition in any particle continues throughout the entire charge or discharge process. The progress of the phase transition is gradual and continuous.

- 2) During discharge, the phase transition starts from the outer surface of the particle and moves inward and a distinctive two-phase interface separates the two end members (LiFePO₄ and FePO₄). The existence of this two-phase interface together with the coexistence of two end members in one particle was experimentally validated by several research groups [80, 81].
- 3) In the process of charge and discharge, a certain level of solid solubility is predicted between the two end members. In other words, end members of composition Li_xFePO₄ and Li_{1-x}FePO₄ (x and x have sm all values close to zero) are predicted.

Although this core-shrinking model gives insight of the process, it didn't necessarily assume all the parameters realistically or in other words it has some inherent limitations. The major limitations include:

- It did not include kinetically induced effects; so it cannot predict the nonequilibrium phase transformation behavior at higher discharge rates.
- It did not include the physics that the Li diffusion through the two-phase interface is significantly sluggish [82].
- It uses isotropic diffusion of Li ions. But Li ion diffusivity has significant different values along different crystallographic directions [51, 83].
- It did not study the particle shape and size effect on the phase transformation behavior.

2.3.3.2 Domino Cascade Model

C Delmas, et.al studied the phase transformation behavior of LiFePO₄ using a model which they call domino-cascade model [73]. In this model they considered the anisotropic Li ion diffusion through the LiFePO₄ lattice where Li ions have a preferentially

high diffusion coefficient along the 'b' direction. They found that once the de-lithiation starts, the change in the charge of Fe ions causes the strong change in the Fe-O and O-O bond lengths in the FeO₆ octahedra which leads to structure distortions. Since FeO₆ octahedron and PO₄ tetrahedron shares a common O-O edge, the change in O-O bond length also affects PO₄. These changes in FeO₆ and PO₄ environments make it difficult to introduce a significant amount of Fe³⁺ ions in LiFePO₄ and a significant amount of Fe²⁺ ions in FePO₄ and hence a large compositional domain for solid solutions is not possible. As a result, a two-phase mixture of two end-members [Li_{1-ε}FePO₄ (ε~0) and Li_ε·FePO₄ (ε'~0)] exists during the charge-discharge process. When the de-intercalation starts a small polaron of Fe³⁺ forms [84] (similar phenomenon happens for Fe²⁺ when lithiation takes place) and the change in the environment of FeO₆ and PO₄ push this polaron to next FeO₆ octahedron. Thus with the lithium de-intercalation along the '*b*' direction a two-phase interface moves along the '*a*' direction which as shown in Fig. 2.8.



Figure 2.8: Schematic view of the 'domino-cascade' mechanism for the lithium deintercalation /intercalation mechanism in LiFePO₄ crystallites [73].

Due to the large difference in bond lengths and angles between LiFePO₄ and newly formed FePO₄ block, strong distortions appear locally that causes instability to the adjacent [FeO₄]_n layers. This results in the easiness to remove lithium ions from the *b* tunnels adjacent to newly formed FePO₄ block instead of starting a new nucleation site to form a new FePO₄ block elsewhere in the same particle. Rather nucleation of a new FePO₄ block within the same crystal particle requires higher activation energy whereas lithium de-intercalation adjacent to the already formed FePO₄ block can proceed without the increase of activation energy. So, once the intercalation or de-intercalation starts in any particle, it finishes very quickly for that particle. But chemical and electrochemical lithiation/de-lithiation can give different pictures. Chemical intercalation/de-intercalation obviously has the tendency to nucleate simultaneously at different points within one crystal particle and hence forms the domains of two phases (LiFePO₄ and FePO₄) in one single crystal [79, 81].

Key attributes of the domino-cascade model are as follows:

- 1) Once the phase transition starts in any particle, it finishes very quickly in that particle. In one particle, phase transition does not continue during the entire period of the discharge process of the cell. Depending on the discharge rate, some particles with good electronic conductivity, having access to the Li ions get activated for phase transition while some other particles might not be activated for transition.
- 2) During the phase transition, only two end members (LiFePO₄ and FePO₄) are present as the unique structure of LiFePO₄ and FePO₄ permits almost no solid solubility, if not, very limited solubility. Once Li ions enter FePO₄ structure (vice versa), distortion in local crystal structure causes
the entire structure transforms into LiFePO₄ very quickly permiting very limited solulibility of Li ions.

2.3.3.3 Comparison between phase transition models

The phase transition models are summarized in the schematics of Fig. 2.9 [85]. The classical core-shrinking model proposed by Padhi et. al [4] and further advanced by Srinivasan and Newman [75] considered isotropic Li ion diffusion (Fig. 2.9(a)). The 'radial core shell model' is a version of core shrinking model but considers anisotropic Li ion diffusion along 'b' direction and in this model it is assumed that during lithiation the Li ion intercalation starts from the surface of the particles and during de-lithiation the Li ion deintercalation starts from the middle of the particles (Fig. 2.9(b)). Domino-cascade model uses the anisotropic Li ion diffusion along 'b' direction, and once transition starts in a particle it very quickly finishes in that particle (Fig. 2.9(c)). Another model, which can be described as a spinodal decomposition model (Fig. 2.9(d)), predicts that a new phase can nucleate at any crystal defects and at any intermediate stage of charge or discharge. It also predicts that one particle can have many domains of both phases and the second phase can be found isolated inside the mother phase. Among these models, only domino cascade model predicts that at any stage of charge or discharge, some particle can be found which has completed the phase transition while some other particles might not even start transforming. Other models predicts the continuous and gradual phase transition progress during the entire period of charge or discharge.

2.3.4 Experimental techniques to study Phase Transition of LiFePO₄

Various techniques have been employed to study the phase transition of LiFePO₄ materials to date. To study the Li ion insertion behavior, different research groups used either chemical [60, 79-81, 86-91] or electrochemical [76, 85, 90, 92-101] lithiation/de-lithiation processes. Besides probing these phase transition was attempted using both



Figure 2.9: Schematic representations of the LiFePO₄ (in red) and FePO₄ (in green) phases in the four models proposed in the literature for the Li insertion/extraction mechanisms in FePO₄ /LiFePO₄ materials: (a) bulk core shell model, [4, 103] (b) radial core shell model, [79] (c) domino-cascade model, [73] and (d) a 4th model that can be called as the spinodal decomposition model freely schematized from ref [80]. In all of the models, except the domino-cascade model, the particles are mixed and composed of the two phases [85].

ex-situ [60, 79-81, 85-89, 91, 93, 94, 97, 99] and in-situ [76, 90, 95, 96, 98, 100-102] methods. Because of these variations of the lithiation/de-lithiation processes and the data recording methods, observation and conclusion about the phase transition mechanism for LiFePO₄ materials varied widely. Various types of characterization tools have been employed by the researchers to probe the phase transition phenomena and to characterize the phase boundary and phases present during the lithiation/de-lithiation process. Among them, the use of Transmission Electron Microscopy (TEM/HRTEM) [81, 85], Raman spectroscopy [80, 91, 102], X-ray Diffraction (XRD) [80, 91, 95, 97, 98, 100,

101], Electron Energy Loss Spectroscopy (EELS/HREELS) [79, 85, 87], Valence Electron Energy Loss Spectroscopy (VEELS) [88], Mossbauer Spectroscopy [86, 100], X-ray Photoelectron Spectroscopy (XPS) [94], Neutron Powder Diffraction (NPD) [60, 76, 86], X-ray Absorption Spectroscopy (XAS) [90, 91, 96, 98, 101], X-ray Absorption Near Edge Structure (XANES) [96], Scanning Transmission Electron Microscopy (STEM) [89], High Angle Annular Dark Field Scanning Transmission Electron Microscopy (HAADF-STEM) [89], Annular Bright Field Scanning Transmission Electron Microscopy (ABFSTEM) [99], Precession Electron Diffraction (PED) [85] etc. has been reported so far. They used active materials of different particle sizes and morphologies. A comparative study on use of different characterization tools is done by Corey T Love et. al. [96]. The key findings from these experimental characterizations can be listed as follows.

- Domains of two end members (Li-rich phase and Li-poor phase) coexist in one particle separated by an interface. This was observed in the exsitu experiments on chemically lithiated/de-lithiated particles [79-81, 86, 87].
- 2) Interface between the two end members are observed. The interface was found because of the existence of a mixture of two end members but not a solid solution [79]. HRTEM also captured a disordered interface (not amorphous) in an ex-situ experiment of chemically de-lithiated large plate shaped LiFePO₄ particle [80, 81]. Using HREELS from the ex-situ experiments, the interface is found to be approximately 8-15 nm wide in a partially chemically de-lithiated plate shaped particles.
- 3) Some theoretical study and models predicted that the interface between the two end members moves along 'a' direction of the orthorhombic crystal structure during the progress of phase transition. HRTEM study in

an ex-situ experiment of partially de-lithiated large plate shaped LiFePO₄ particle [81] and XPS study [94] on samples of electrochemically lithiated/de-lithiated at different state of charge-discharge also were able to capture this phenomenon. One dimensional movement of interface was also recorded using kinetic study and analyzed with the Avrami-Johnson-Mehl-Errofev equation [92, 104].

- 4) Domino-cascade model [73] predicted at any stage of charge or discharge, electrode consists of a mixture of fully lithiated particles and completely de-lithiated particles. In one particle two end members does not coexist. This is also supported by ex-situ study on electrochemically de-lithiated LiFePO₄ particles using HRTEM and PED [85].
- 5) Two end members were identified as Li_{0.89}FePO₄ and Li_{0.05}FePO₄ at room temperature from ex-situ neutron diffraction data from approximately 100 nm particles [60]. Smaller particles have higher solid solubility [92, 105] at room temperature and can even be one single phase of composition LixFePO₄ (x can have any value within 0 and 1) for less than 15 nm particles [105] for the entire period of charge/discharge process. In other words, the miscibility gap between the two end members during charge/discharge process can completely disappear at high temperatures for LiFePO₄ particles below certain critical size.
- 6) Usually the LIBs containing LiFePO₄ material are charged/discharged within the potential window of 2.5 V to 4.2 V. But if it is deeply discharged to 0.75 V or so, below 1.5 V the battery gives a slopping (reducing or increasing during discharge and charge respectively) potential profile

which indicates the absence of any two-phase system and one solid soluble phase [16].

2.4 Scope of study and objectives

It is already discussed that phase transition between LiFePO₄ and FePO₄ gives a flat plateau of potential during charge or discharge and at the particle level understanding of the detailed mechanism of this phase transition is very important to overcome the high rate charge/discharge barrier and also to optimize the battery design for various applications including hybrid electric vehicles. So far the observation of phase transition phenomena is contradicted by different phase transition models and also by experimental observation using a range of different characterization tools and experimentation methods. Based on all the information presented here, the phase transition of LiFePO₄ material needs to be studied in more detail and hence the objectives of this study are listed as follows.

- Synthesize highly crystalline LiFePO₄ materials and characterize them using different tools such as particle size analyzer, XRD, SEM, electrochemical tests, etc.
- 2) Since chemical lithiation/de-lithiation and electrochemical lithiation/delithiation gives different observation and in a real working battery electrochemical lithiation/de-lithiation takes place, it is critical to study LiFePO₄ materials through in situ electrochemical test.
- In-situ experiment gives real time information from a working battery and in-situ study is always preferred over ex-situ experiments; therefore insitu techniques were adopted herein to study the phase transition of LiFePO₄.

- 4) Phase transition of LiFePO₄ needs to be studied at different length scale such as at particle lever and at the bulk level, and hence the in-situ experiments were conducted at two distinctly different length scales. For example, in-situ XRD experiment in a traditional powder X-ray diffractometer can give the bulk-level phase transition information of LiFePO₄ but Raman spectroscopy could be used to probe very small electrode area so that particle-level phenomena can be studied.
- 5) To better understand the phase transition of LiFePO₄ at the particle level, numerical study will be conducted to understand the effect of particle size and particle morphology on phase transition.

Chapter 3

Synthesis of LiFePO4 Particles and Li ion Battery Making

3.1 Introduction

Both synthesis of LiFePO₄ materials and LIB fabrication process need special care as it requires the special inert atmosphere. Electrode making process has to be tuned properly to obtain optimum electrode properties like porosity, conductivity, etc. and the battery has to be assembled in an argon filled glove box to protect from atmospheric oxygen and moisture. Special attention has to be paid to all these requirements to make a good performing LIB. In this chapter, synthesis of active LiFePO₄ material, their characterization, electrode fabrication using them, and finally battery assembly process is discussed in detail.

3.2 Synthesis of LiFePO₄

3.2.1 Solid State Synthesis

LiFePO₄ material has been synthesized using solid state synthesis method. In the beginning of the solid-state synthesis, the precursors such as Li_2CO_3 (Sigma-Aldrich), FeC₂O₄.2H₂O (Sigma-Aldrich) and NH₄H₂PO₄ (Sigma-Aldrich) were mixed properly in a stainless steel container using ball milling for 1 hour. The precursor mixture was heattreated at 350°C in argon for 3 hours in a tube furnace. During the heat treatment, the precursor compounds decomposed and released H₂O, CO₂ and NH₃. After this step, the carbon coating precursor cellulose acetate (Sigma-Aldrich) was mixed with the decomposed mixture and was heated in a tube furnace at 700°C for 8 hours under argon gas flow to get carbon coated crystalline LiFePO₄. The solid synthesis method is shown with the help of the flow diagram in Fig. 3.1.



Figure 3.1: Flow diagram for solid state synthesis method to synthesize LiFePO₄

material.

3.2.2 Characterization

After synthesis, the routine characterization techniques like XRD, SEM, electrochemical test, etc. have been employed to study the material. X-ray diffraction test was done to confirm the phases of the synthesized material. Figure 3.2 shows the XRD patterns of the synthesized LiFePO₄ and the olivine structure (*Pnma*) was confirmed using the JCPDS card number 81-1173.

Particle size of the synthesized LiFePO₄ particles were determined using Nanosight Nanoparticle Tracking Analyzer and the result is presented in Fig. 3.3. The average particle size determined was approximately 222 nm while most of the particles were approximately 196 nm with a standard deviation of 74 nm.



Figure 3.2: XRD patterns obtained from materials synthesized through solid state synthesis method. JCPDS card # 81-1173 is shown in bottom plot.

SEM image of the particles also give an indication of the particle size and also the particle morphology (Fig. 3.4). SEM study of the LiFePO₄ material was done using ZEISS Supra 55 Variable Pressure FE-SEM. Particles obtained through solid state synthesis route have an irregular shape as seen in SEM image (Fig. 3.4).



Figure 3.3: Particle size distribution of the synthesized LiFePO4 material at 700 $^\circ$ C for 8

hours.



Figure 3.4: FE-SEM image of the solid state synthesized LiFePO₄ material at 700 °C for 8

hours.

3.3 Electrode Fabrication

To make the cathode, the first step is to make slurry containing the synthesized LiFePO₄ and additional conductive carbon. The steps of making electrode slurry are shown in Fig. 3.5. First, a glue was prepared e when binder PVDF (Polyvinylidene fluoride, Kynar HSV900, Elf Atochem) was dissolved in NMP (1- Methyl, 2-Pyrrolidinone, Acros Organics) at 50 °C. Then a mixture of carbon coated LiFePO₄ (C-LiFePO₄) and additional conductive carbon (Timcal Super P Li carbon) were slowly added to the glue under mixing using a magnetic stir bar on a hot-plate. After the whole mixture is added to the glue the rotation speed was adjusted to the 800 rpm and the mixing continued for 12 hours at 50 °C. The composition of cathode mixture varies in the study; one typical composition was 70% C-LiFePO₄, 20% conductive carbon, 10% PVDF binder. The amount of liquid NMP is also important and it has to be optimized for particular conductive carbon to get the right viscosity for coating on current collectors using a doctor blade. If the surface area of conductive carbon is high, more liquid NMP is required.



Figure 3.5: Schematic of slurry making process containing active material and conducting carbon.

By the time the slurry is being prepared, the current collector AI foil has to be ready to be coated by active material slurry. A piece of highly pure AI foil (Li ion battery grade) was cut and attached on a flat glass plate. The surface of the AI foil that will be coated by the active material slurry was pre-coated with a thin coat of graphite by rubbing a graphite bar on AI foil. This graphite layer later on acts as a solid lubricant during pressing the electrode after drying. Without this graphite layer, once the coating of active material is pressed to obtain a porous but dense continuous conductive network of





carbon, the coating tends to be peeled off from the Al current collector. After applying the thin graphite layer, the mixed C-LiFePO₄ slurry was poured on the graphitized area of the Al foil with the help of a doctor blade while the slurry was coated. The blade gap used typically for the experiments is approximately 350 μ m and after drying and pressing the active material coating thickness became approximately 50 μ m. The steps of slurry coating are shown in Fig. 3.6. After the coating was done, it was dried in an oven at 80 °C for 30 minutes. The cathode was then dried in a vacuum oven at 120 °C for 12 hours after the primary evaporation of liquid NMP at 80 °C. An appropriate size of the cathode

was then cut and pressed using 2 ton/cm² pressure before being used as a cathode. Cross sectional SEM image of a dried and pressed cathode is shown in Fig. 3.7.



Figure 3.7: SEM image of a cathode containing 70% C-LiFePO₄, 20% Conductive carbon and 10% PVDF binder.

Anode used for the battery for later experiments is lithium ribbon. Lithium ribbon was cut and cleaned to remove the thin oxide layer by scratching them off and exposing fresh lithium metal surface inside an argon filled globe box. The cleaned lithium metal ribbon was then pressed against a precut stainless steel mesh. As the lithium metal is very soft, the steel mesh penetrated inside the lithium with pressure which established a very good electrical connection. This steel mesh was connected to an external potentiostat which was then electrically connected to the lithium anode.

3.4 Electrolyte Preparation

A 1:1 volume ratio of EC:DMC (Ethylene Carbonate, Di-methyl Carbonate) solvent system was used to prepare the electrolyte. EC is in a solid at room temperature and hence was melted and taken inside the argon filled glove box to prepare the electrolyte. The salt LiPF₆ which was dissolved in the solvent system is very sensitive to moisture and hence was only handled inside the glove box. After opening the LiPF₆ bottle, the solvent mixture was added immediately to minimize the exposure of LiPF6 to the atmosphere even though the glove box has very little moisture. 1M LiPF₆ solution was prepared to be used as electrolyte in all subsequent battery assembly processes.

3.5 Battery Assembly

For battery assembly, a home-built cell[106] was used which has a sealed glass cavity to house the cathode-separator-anode sandwich and electrolyte, with a sealant cap made of epoxy. All the battery assembly procedures were performed inside the argon filled glove box and also all the items including the cathode, separator, etc. that were taken inside the glove box were dried for at least 6 hours to ensure that they are moisture free. First, the cathode-separator-anode sandwich was inserted inside the glass cavity where the Celgard® 2500 separator was used and tested for any possible short circuit. If there was no shorting, the cathode and anode were connected with terminals inside the glass cavity which were then connected to the external stainless steel wires. Two small PTFE (poly tetrafluoro ethylene) wedges were used to push the cathode-separator-anode sandwich against the glass wall to ensure the least distance between the electrodes and hence to reduce the ohmic potential drop. Electrolyte was then added to the glass cavity to soak the separator and electrodes. Finally, the sealant cap was attached to protect the battery assembly. As the cell had the transparent glass housing, the battery was then

wrapped with some opaque Al foils to protect it from light which might degrade the electrolyte.

The assembled battery was kept idle for few hours and then subjected to the necessary electrochemical tests like cyclic voltametry (CV), impedance test, and most importantly galvanostatic charge/discharge tests. One such typical galvanostatic charge and discharge performance of a cell having a cathode of 70% C-LiFePO₄, 20% Conductive carbon and 10% PVDF binder is shown in Fig. 3.8. The second and third cycle of charge and discharge profile is presented in Fig. 3.8 where the columbic efficiencies are 93.5% and 93.9% which indicates the stability of the battery.



Figure 3.8: Galvanostatic charge discharge profile of a lithium ion cell. Cathode Composition is 70% C-LiFePO₄, 20% Conductive carbon and 10% PVDF binder. C_2 and D_2 indicate the 2nd cycle charge and discharge profile while C_3 and D_3 indicate the 3rd cycle charge and discharge profile.

Chapter 4

In-situ Studies using Raman Spectroscopy

4.1 Introduction

Although LiFePO₄ has so many attractive properties like stable potential during discharge, safe, and less expensive, the fundamental mechanism of Li transport during charge-discharge and the phase transition between two end members in LiFePO₄ is not fully understood. As has been already discussed in the previous chapters, two issues are controversial regarding the phase transition of LiFePO₄ material: one is the progress kinetics of the phase transition within a particle and the other is the existence of any possible solid solution with composition of Li_{1-x}FePO₄. So far, many phase transition theories/models have been proposed and many of them are supported by different experimental data as already discussed in Chapter 2. The different experimental techniques that are either ex-situ or in-situ, and the different type of intercalation/deintercalation processes of Li ion that are either electrochemical or chemical lead to varied observations. Although in-situ experimentations have many practical challenges to obtain meaningful results with minimum errors, they are very crucial to get the real time phase transition information. At the same time, electrochemical intercalation/de-intercalation studies provided more comprehensive understanding of a real operating battery than chemical intercalation/de-intercalation study.

Raman spectroscopic technique has been used to study different battery components like electrolyte [107-110], carbon additive [111-114], carbon coating on the active material particles [115, 116], active material LiFePO₄ [117, 118], and other electrode components [37, 44, 113, 117, 119, 120] as well. If electrochemical charge/discharge test can be performed together with Raman spectroscopy, the result

can be more meaningful as the results come from a real working battery. So in this work the in-situ Raman spectroscopy was performed while a battery is being charged/discharged to study the phase transition phenomena of LiFePO₄ material. As the Raman laser spot size is around 2µm, it can practically study particle-level phase transition process. In this chapter, the detailed experimentation technique for in-situ Raman spectroscopy of a LIB containing LiFePO₄ material and the associated phase transition results during different rate charge and discharge process have been discussed.

4.2 Design of Battery Casing

LIBs require a sealed container to house the components so that they are protected from the atmospheric moisture and oxygen as has already been mentioned in Chapter 3. To assemble a battery to be used in the in-situ Raman spectroscopy, additional design requirements have to be considered to allow the Raman laser access the active material of the electrode to be studied. To do this, a Raman laser transparent window has to be installed at the positive electrode side of the battery.

Figure 4.1 shows the design of the transparent cell housing and experimental setup for the in-situ Raman study. The cell housing was made of an airtight cavity of glass as the electrolyte reservoir, an epoxy connector, and a sealing cap. The conductive terminals of the cell were made of stainless steel for both the cathode and anode.

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Figure 4.1: Schematic of the cell design for in-situ Raman spectroscopy. The real working cell operating in the Raman spectrometer sample holder is shown in the bottom.

This cell configuration prevents the cathode surface to be covered with an excessive amount of electrolyte. This feature consequently facilitates deeper penetration of the laser beam and therefore improves detection of chemical changes on the cathode surface. A NaCl single crystal (Edmund Optics) used as the transparent window for the

laser pathway not only minimizes laser intensity loss but also has well-defined Raman signatures that do not overlap with signals from the battery materials.

4.3 Electrode Fabrication and Battery Making

4.3.1 Electrode Fabrication

Electrodes for in-situ Raman study are different comparing to the electrodes that are used in the normal Li-ion batteries. The positive electrode in normal batteries has active material (LiFePO₄) coated on the current collector (Al foil) and the side with active material coated faces towards the separator and the negative electrode (anode) in a working battery. But with this configuration it is not possible to study using the in-situ Raman spectroscopy because the current collector (aluminum foil) blocks the active materials from being directly accessed by the Raman laser. So, the electrodes that will be used in the batteries for in-situ Raman study have active materials coated on a homemade aluminum mesh where small holes of the aluminum mesh will allow the Raman laser to interact with the active material. The hole size on aluminum mesh is approximately 0.5 mm and the Raman laser spot size is 2 µm which makes possible to study the material. And also, as the individual hole can be physically marked, it is possible to study the same spot if the experiment is discontinued. The preparation of this special kind of electrode is schematically shown in Fig. 4.2.



Figure 4.2: Schematic of the cathode making process specially suitable for in-situ Raman

spectroscopy.

Active material slurry contains 90% C-LiFePO₄, 2% conductive carbon, 8% binder (PVDF). This slurry preparation is similar to the process described in Chapter 3. The slurry was then coated on the aluminum mesh (Fig. 4.2), dried in a vacuum oven at 120 °C and pressed with a 2 ton/cm2 pressure before being used in the in-situ Raman study. Finally, the electrode was cut to a size of 1.5 cm × 1.5 cm. Since the Raman laser interacts with the carbon (either carbon additives or carbon coating on LiFePO₄) in the electrode, some corrective measures were taken which will be described in section 4.3.4.

For the anode, an appropriate size (1.5 cm \times 1.5 cm) of lithium ribbon was cut and pressed against a steel mesh, which acted as a current connector.

4.3.2 Battery Making

Positive electrode, separator, and the battery components were dried for 6 hours and then transferred to the argon filled glovebox to assemble the battery for the in-situ Raman study. 1M LiPF₆ solution with a solvent mixture of 1:1 (volume ratio) of DMC:EC was prepared inside the glovebox to be used as electrolyte for the battery. The anodic lithium, Celgard 2500 separator, and the specially fabricated cathode were assembled as a sandwich which was subsequently inserted inside the battery housing (Fig. 4.1) and then with the help of a multi-meter any possible short shorting was examined. If there was no shorting, a PTFE wedge was pushed beneath the anode side to bring the cathode closer to the Raman laser transparent window. Two electrodes were then attached to the external electrical connection and finally electrolyte was added to the electrode sandwich. After the electrodes and separator were soaked by the electrolyte for 30 minutes, extra electrolyte was drained to enhance the Raman signal from the active material. The sealant cap was finally attached properly before the battery was taken outside the glovebox for the in-situ Raman spectroscopy.

4.3.3 Raman Laser Interaction with battery components and Corrective measures

The purpose of this in-situ Raman study is to understand the phase transformation of LiFePO₄ materials. But to get the best results, Raman laser interaction with all the battery components has to be considered. Table 4.1 lists all possible Raman active vibration modes of LiFePO₄ materials and its de-lithiated phase FePO₄, and also other vibration modes for different battery components such as carbon and electrolyte solvents. The positive electrode was composed of carbon coated LiFePO₄, additive carbon, and PVDF binder. Both of these coated carbon and carbon additives along with the electrolyte interact with Raman laser and complicate the study of LiFePO₄ phase transition.

Table 4.1: List of Raman active vibration modes for components (electrode material and electrolyte) of the working battery.

Material	Vibration Mode	Peak Position,	Reference
		cm ₋₁	
LiFePO ₄	Asymmetric stretching of PO_4^{3-} (u3)	1067, 995	[102]
	Symmetric stretching of PO ₄ ³⁻ (u1)	949, 943	[102]
FePO ₄	Asymmetric stretching of PO_4^{3-} (u3)	1126, 1068	[102]
	Symmetric stretching of PO ₄ ³⁻ (u1)	959	[102]
	Symmetric bending of PO ₄ ³⁻ (u2)	665	[102]
	Asymmetric stretching of PO_4^{3-} (u4)	588	[102]
γ-Li ₃ Fe ₂ (PO ₄) ₃	PO_4^{3-} deformation	448, 586	[117]
	Internal stretching of PO ₄ ³⁻	990, 1040,	[117]
		1119	
α-Fe ₂ O ₃	O-Fe-O bending	395	[117]

Table 4.1—Continued

	Fe-O stretching	600	[117]
	Disorder band due to point defects	646	[117]
	Second order disorder band due to point defects	1280	[117]
С	C-C in-plane bond stretching mode E_{2g} of sp ² bonded graphite (G band)	1585 or 1582	[114, 121]
	Breathing mode of sp ² sites A _{1g} in hexagonal rings of carbon in graphite planes (D band)	1350	[122]
	Disordered sp ³	1194, 1490, 1620	[115, 121, 123, 124]
EC	Symmetric stretching of C–O	900, 905 (salt solvated)	[125]
	Bending of C=O	720, 730 (salt solvated)	[125]
DMC	Symmetric stretching of C–O	920, 930 (salt solvated)	[125]

Some corrective measures were taken to understand these Raman interactions to get the phase transition signals from the active material LiFePO₄. LiFePO₄ particles were coated with carbon during the synthesis to enhance the electronic conductivity, but this coated carbon significantly diminishes the Raman vibration signal from the LiFePO₄ [117]. The coated carbon has both sp² bonded graphitic carbon (G band) whose position is at around 1580 cm⁻¹, and also disordered sp³ bonded diamond carbon (D band) whose position is at around 1340 cm⁻¹. The position of G band and D band is marked in Fig.

4.3D. This sp² bonded graphitized carbon is beneficial because it enhances the electronic conductivity to the LiFePO₄ particle, particularly at the surface where electrochemical reaction takes place, but it gives hindrance to the optical penetration of the Raman laser. The Raman laser skin depth δ is expressed as

where ω is the optical frequency, μ is the permeability, *c* is the speed of light, and σ is the electrical conductivity of the sample. Once the carbon coating from the LiFePO₄ is removed, the conductivity is reduced causing the laser optical skin depth increased to get the enhanced signal from LiFePO₄. So, carbon coating from the LiFePO₄ particles has to be removed partially from the areas to be studied. This ensured minimal detrimental impact on the phase transformation because most of the areas on the electrode were unaltered to make a successful working battery for the in-situ Raman spectroscopy. Raman laser simply could gasify this coated carbon without altering the LiFePO₄ in olivine structure [117]. Before testing the battery the cathode was exposed to 10 mW Raman laser for 5 minutes to gasify the coated carbon to carbon dioxide. But during gasification of coated carbon, the decomposition of LiFePO₄ can also happen according to the following reaction.

$$12\text{LiFePO}_4 + 3\text{O}_2 \rightarrow 4\text{Li}_3\text{Fe}_2(\text{PO}_4)_3 + 2\text{Fe}_2\text{O}_3$$



Figure 4.3: Raman spectra from four different areas after carbon gasification treatment. Areas A and B show decomposition of LiFePO₄ while areas C and D shows the vibration signal of PO_4^{3-} (u1) of LiFePO₄ material [106].

Special care has been taken to minimize or totally avoid LiFePO₄ decomposition by the Raman laser. Fig. 4.3 [106] shows Raman spectra of four different areas (150 µm × 150 µm) after the gasification of coated carbon on LiFePO₄ particles. Areas C and D in Fig. 4.3 show the presence of vibration mode (u1) of LiFePO₄ while areas A and B do not show the Raman signal from LiFePO₄, but rather some other vibration signals which come from the decomposition product of γ -Li₃Fe₂(PO₄)₃ at 488, 990 and 1040 cm⁻¹ [117] and from the decomposition product of Fe₂O₃ at 217, 277, 395, and 600-650 cm⁻¹ [117]. So, the C or D area was considered for in-situ study area as they give the LiFePO₄ signature signal and did not show any decomposition of LiFePO₄.

4.4 In-situ Raman experiment: Data Recording and Analysis

In Table 4.1, it has been already shown that there are several Raman active vibration modes for both LiFePO₄ and FePO₄. Among them, Raman bands of LiFePO₄ (u1) at around 950 cm⁻¹ and of FePO₄ (u1) at around 960 cm⁻¹ were targeted for study using in-situ Raman spectroscopy as these signals are more intense than others. After making the battery and before the in-situ study, the areas C and D (Fig. 4.3) were mapped for the intensity of the LiFePO₄ signal (950 cm⁻¹) to find the suitable spot for further in-situ investigation. Such a mapping ($80 \ \mu m \times 80 \ \mu m$) for signal at 950 cm⁻¹ is shown in Fig. 4.4 [106]. The spots with intense LiFePO₄ signal (950 cm⁻¹) were considered for further in-situ study.



Figure 4.4: Raman mapping of a area after carbon gasification process of the electrode. The points with intense signal from the vibrational signal of PO_4^{3-} (u1) of LiFePO₄ material can easily be found from this map which can be used for subsequent in-situ study [106].

In-situ Raman studies were conducted using a Thermo Scientific DXR with a 780 nm laser. The laser spot size was 2 μ m with a laser power of 2 mW using 10X magnification objective lens. The exposure time was 16 s with a photo-bleach time of 30 s. The schematic in Fig. 4.5 shows the microscopic description of the electrode under the

in-situ Raman laser. The average particle size used in this experiment was around 200 nm as is shown in Chapter 3; the laser spot size was 2 µm and the Raman laser focused area contained only a few active material particles. Thus the in-situ Raman study would be able probe the particle-level phase transition phenomena.



Figure 4.5: Schematic showing electrode configuration for in-situ study. The small laser spot only focuses on few active material particles.

4.4.1 First Cycle: Slow discharge

During the in-situ Raman study, the Li ion cell was electrochemically charged and discharged using a multi-channel battery tester. The cell was charged and discharged for three cycles; each cycle was at different rates. The first cycle started with a constant-current charge at 0.3C until the cell voltage reached 4.0 V, then followed by a constant-voltage charge at 4.0 V for 20 minutes. After that rest period of 45 minutes was introduced and the cell was then discharged at 0.3C rate until the voltage reached 2.5 V,

then followed by a 1 hour rest period. The first-cycle electrochemical test result is presented in Fig. 4.6 where the black vertical lines on the potential profile curve indicate the points where the Raman spectra were collected.



Figure 4.6: Charge discharge cycle at 0.3C rate during slow rate in-situ Raman spectroscopy. Blue line indicates the current profile and the black line indicate the potential profile and vertical lines on them indicates the point when the Raman spectra were recorded.

The associated Raman spectra that were collected during the 1st cycle of the electrochemical test are presented in Fig. 4.7. From the spectra (Fig. 4.7) it is clearly observed that during charge the LiFePO₄ particles transformed suddenly to FePO₄ between scan C5 and C6 (Figure 4.7). The same feature was observed during the discharge process where at the very beginning the signature LiFePO₄ phase peak showed up from scan D2, indicating that all the particles under the laser spot got transformed completely. To understand it better, quantitative analysis was performed on each spectrum by doing the Gaussian fit analysis. The Gaussian fit analysis of a representative spectrum (scan D1, 1st cycle) that contains both distinctive signature peaks of LiFePO₄ and FePO₄ is shown in Fig. 4.8.

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Figure 4.7: Raman spectra collected during electrochemical charge discharge at
0.3C. R1, R2, and R3 indicates the Raman spectra collected during rest period and
C1, C2, etc indicates the Raman spectra taken during charge and D1, D2, etc are
Raman spectra taken during discharge as indicated in Fig. 4.6.

During the Gaussian fit analysis of any spectrum, constant baseline correction was applied and the same level of correction was applied to all of the spectra to keep the quantitative analysis unbiased. From the Gaussian fit analysis of each spectrum, the peak intensity information of each peak (LiFePO₄ peak and FePO₄ peak) was taken to calculate the relative amount of phases present at any point of charge or discharge. For example, as shown in Fig. 4.8, the intensities of LiFePO₄ and FePO₄ peaks are 1.53 and

1.66, respectively, corresponding to the relative amounts of 48.03% and 51.97% for the $LiFePO_4$ and $FePO_4$, respectively, at that point.

The relative amount of phases during the charge and discharge process obtained from the Gaussian fit analyses are plotted and shown in Fig. 4.9. It is observed that during the charge process the particles in the probed area under the Raman laser transformed abruptly between scan C5 and C6; and during the discharge process, the study area of the electrode finished phase transition within the first two scans. This indicates that once phase transition starts within that small area of the electrode, it finishes very quickly.



Figure 4.8: Peak de-convolution and Gaussian fit analysis of one representative Raman spectra D1 collected during electrochemical discharge at 0.3C. The time when D1 is indicated by the 1st vertical line on the potential profile of discharge curve in Fig. 4.6.

It is already explained that the in-situ Raman study only investigated a few particles under the laser focus. Though the electrode microstructure is very much heterogeneous considering a large area of electrode, the small local area (Raman laser spot size of 2 µm) is quite homogeneous in terms of the local electronic conductivity and the availability of Li ion through the porous network of electrode. So, the particles in a small local area complete phase transition almost together and hence once phase transition starts it gets completed very quickly. Scan D1 shows that half of the particles got transformed within the probed area, but the very next scan indicates the completion of phase transition for all the probed particles. This clearly shows that in any individual particle once the phase transition starts, it completes very quickly. This particle-level phenomenon is in good agreement with the phase transition mechanism as explained by the domino cascade model [73].



Figure 4.9: (a) Progress of phase transition during charge and discharge at 0.3C.(b) Current and Potential profile during charge and discharge is shown along with the points when Raman spectra were recorded.

4.4.2 Effect of Discharge rate: Second and Third cycle

The in-situ Raman study was also performed for two more cycles. In the second cycle, the cell was electrochemically charged at 1C rate until the cell voltage reached 4.0 V, then followed by a constant-voltage charge for 50 minutes at 4.0 V. After that a rest period of 35 minutes was introduced. The cell was then discharged at a constant rate of 1C until the voltage reached 2.0 V, subsequently followed by a rest period of 30 minutes. The electrochemical charge-discharge test results are presented in Fig. 4.10(a) and the associated Raman spectra are presented in Fig. 4.10(b).

At a higher rate charge such as 1C, it is clearly observed from Fig. 4.10(b) that the phase transition was delayed and in some particles the phase transition did not take place at all. The last scan of Raman spectrum (C5) during the constant-current charging shows a small shoulder of the FePO₄ phase, and during the constant-voltage charging almost half of the particles in probed area got transformed even though the charging current was reduced as seen from Fig. 4.10(b). Finally at the end of the charging process, approximately half of the particles got transformed.

During the discharge at 1C rate, the first scan D1 shows almost no transformation from the rest cycle 'R2'. But the second scan D2 shows that most of the particles (not all) that transformed during the charging process transformed back to LiFePO₄. Only a weak FePO₄ peak remained and that weak FePO₄ peak is present in all the scans after D2. So this high-rate charge and discharge show that even the Raman laser spot was very small (2 μ m) and it contains only a few particles, heterogeneity arises at high rate at the very local area. Some particles within the perimeter of a few μ m do not transform simultaneously, leading to an obvious delay in phase transition at a higher rate.

The third cycle of the electrochemical test was performed at a even higher rate of 2C. The cycle started with a constant-current charge at 2C rate until the potential reached

4.0 V, which was then followed by a constant-voltage charge at 4.0 V for 15 minutes and then by a rest period of 15 minutes. The cell was then discharged at 2C rate until the voltage reached 2.0 V; after that the cell was rested for 15 minutes. The electrochemical result for this cycle is presented in Fig. 4.11(a) and the associated Raman spectra are presented in Fig. 4.11(b). From the spectra, it is observed that almost no phase transition took place for the LiFePO₄ particles under the Raman laser as a similar weak FePO₄ peak was observed throughout the cycle. At this high charge/discharge rate, the phase transformation did not take place in any of the particles that were present in the probed area even though phase transformation was observed at a lower rate (0.3V).



Figure 4.10: (a) Progress of phase transition during charge and discharge at 1C. (b) Current and Potential profile during charge and discharge is shown along with the points when Raman spectra were recorded.



Figure 4.11: (a) Progress of phase transition during charge and discharge at 2C. (b) Current and Potential profile during charge and discharge is shown along with the points when Raman spectra were recorded.

4.5 Conclusion

This chapter discusses the detailed experimentation technique and results for the in-situ Raman study. The Li ion cell was charged and discharged at three different rates including 0.3C, 1C and 2C. The Raman laser spot size was 2 µm and the average particle size of LiFePO₄ was about 200 nm, indicating the probed electrode area contained only a few active material particles. During the slow discharge, it was observed that during charge and discharge once the phase transition started it finished very quickly. This finding is in good agreement with the domino-cascade model. During the fast discharge or charge at 1C, though the similar trend of quick phase transition upon being initiated still holds, a general phase transition delay was observed. At the end of charge or discharge process, not all the particles completely transformed. Further increasing the charge and discharge rate causes a very unfavorable situation where almost no phase transition was detected by the in-situ Raman spectroscopy.

Chapter 5

In-situ X-ray Diffraction (XRD) Study

5.1 Introduction

Among all the spectroscopic techniques available, X-ray diffraction is very suitable to study the phase transition of LiFePO₄. However, there is a limitation of using the traditional diffractometer that uses powder or polycrystalline samples: the X-ray irradiated area is quite big and it gives the average results of transformed/untransformed phases. This restricts the opportunity to study the related process at the particle level. In the previous chapter, in-situ Raman spectroscopy has been successfully applied to get the particle-level phase transition information.

In this chapter, the use of an in-situ XRD spectroscopic techniques will be discussed which is targeted to understand the phase transition phenomena of LiFePO₄ material. The chapter will cover details of the in-situ XRD experiment including the design of the battery housing containing X-ray transparent window, electrode preparation that enables investigation of the active material in the reflection mode of X-ray diffraction, electrochemical test protocols, and finally results of the phase transition study.

5.2 Making battery housing for in-situ XRD study

Making a battery seems very easy as it consists of very simple components like electrodes, separator and electrolyte. But Li-ion batteries use non-aqueous electrolyte and the atmospheric moisture when absorbed can degrade the electrolyte thus causing performance/capacity loss. Besides, for our experimental Li-ion batteries fresh lithium metal ribbons were used as the anode which even cannot tolerate trace amount of oxygen and moisture. All these conditions require a sealed housing which will hold the
battery components such as electrodes, separator, and non-aqueous liquid electrolyte. In case of studying the battery with the in-situ XRD the housing has another set of extra requirements which include the compatibility that it needs to fit into the sample holder of the XRD machine, and incorporation of an X-ray transparent window which will allow the X-ray to interact with the electrode materials and the diffracted beams to be recorded. To meet all these requirements, a specially designed battery housing was fabricated in the lab. The schematic design and the working case are shown in Fig. 5.1. Inside the housing near the cap two electrical connectors were placed to conduct current produced from the battery to an external battery tester.



Figure 5.1: Schematic of the cell design for in-situ XRD study is shown on the left side and the real working cell operating in the X-ray diffractometer sample holder is shown on the right side.

As shown in Fig. 5.1, the battery housing is made of epoxy and the X-ray transparent window is made of a Kapton film. The Kapton film has been tested for X-ray

transparency and the XRD pattern of the LiFePO₄ powders under the Kapton film is shown in Fig. 5.2.



Figure 5.2: XRD spectra shown under the Kapton film indicates the suitability of using as a X-ray transparent window material in the electrochemical cell to be used in in-situ XRD study.

5.3 Electrode preparation and battery assembly for in-situ XRD study

Electrodes for the in-situ study are different in design comparing to the electrodes that were used in the normal Li-ion batteries. The positive electrode in normal batteries has active material (LiFePO₄) coated on one side of the current collector (Al foil), which is then facing towards the negative electrode (anode). But with this configuration it is impossible to study the active materials using the traditional X-ray diffractometer.

Traditionally the reflection mode of diffraction is used in XRD studies and the aluminum foil as the current collector blocks the active materials from the direct access of X-ray. So, the electrodes that will be used for the in-situ XRD study have active materials coated on both sides of a homemade aluminum mesh. As the current collector is a mesh instead of a solid film, electrolyte can always percolate to the other side through porous network of active materials and eventually electrochemical reactions can take place throughout the active materials of both sides. One side of the electrode faced toward the anode and the other side was studied using reflection-mode XRD.

The meshed aluminum foil was coated with active material mixture (65% active material, 20% C, 15% PVDF) on both sides using a doctor blade and then dried in a vacuum oven at 120 °C for 5 hours. Subsequently the active material and current collector sandwich was pressed with a 2 ton/cm² pressure. Finally, the electrode was cut to appropriate size (1.5 cm × 1.5 cm) to be used in the battery.

For the anode, an appropriate size $(1.5 \text{ cm} \times 1.5 \text{ cm})$ of lithium ribbon was cut and the oxide layer was scratched off to expose the fresh lithium surface in an argon filled glove-box. This lithium ribbon was then pressed against a steel mesh and this steel mesh acted as an electrical connector with anodic lithium.

Before making a battery, the previously prepared cathode was dried for at least 6 hours and other components such as battery housing, the anode, screws, separator and accessories to make the battery were cleaned and dried for 30 minutes to ensure that they were free from moisture. To assemble the battery the Celgard® 2500 separator was placed between the two electrodes to form a sandwich which was then placed inside the specially designed battery housing (Fig. 5.1). Two PTFE wedges were pushed beneath the electrode sandwich upwards so that the cathode surface reached close to the X-ray transparent window. The next step is to add a freshly prepared electrolyte which was a

1M LiPF₆ solution with a mixed solvent system of 1:1 (volume ratio) of DMC:EC. The final step was to cap the battery housing to ensure proper sealing. The battery was kept inside the glove-box for a few hours so that separator and electrodes were soaked properly and the total system stabilized. After that, the extra electrolyte was drained to avoid the X-ray scattering during in-situ XRD experiment.

5.4 In-situ XRD experimentation

Before placing the battery inside the sample holder of the X-ray diffractometer the battery was tested with a few cycles of charging and discharging. During the in-situ XRD experiment, the electrochemical charging and discharging were conducted using a MTI multi-channel battery tester and the X-ray diffraction test was conducted using Siemens D500 X-ray diffractometer.

5.4.1 Test Protocol and Recording the real-time phase transformation

For the in-situ XRD experiment, the battery was tested at four different rates. Each cycle of electrochemical charging and discharging consisted of five consecutive steps. The first step was to charge the battery with a constant current at 0.1C rate followed by the next step of constant-voltage charge for 15 minutes. The battery was then kept at rest (third step) for 15 minutes. The battery was subsequently discharged at a constant rate (fourth step) followed by another rest period (fifth step) of 15 minutes. During the constant-current charging step, the cell voltage was set to a window between 2.5 V and 4.1 V and the constant voltage charging is maintained at 4.1 V. The battery was discharged at four different discharge rates of 0.1C, 0.3C, 0.5C and 1C. The test protocol of this in-situ XRD study is summarized in Table. 5.1.

The XRD pattern of the synthesized material is already presented in Fig. 3.2. Among all the peaks, LiFePO₄ (311) is targeted for study as intensity of this peak is the

highest and noises from the battery components such as electrolyte block some of the peaks. The position of the LiFePO₄ (311) peak is at ~35.66° (20), and the position of the FePO₄ (311) is calculated as 36.703° considering the lattice parameters of a, b, c as 9.81 A°, 5.79 A° and 4.78 A° [126, 127]. There is one more peak corresponding to LiFePO₄ (121) at 36.6° between LiFePO₄ (311) and FePO₄ (311). Therefore, the X-ray scan range was from 35.1° to 37.1° (20) with 0.01° step size. The X-ray collection time for each step was 1 s and hence the total scan time over the targeted angle range is approximately 200 s. The electrochemical charge-discharge cycle is shown in Fig. 5.3 and corresponding insitu XRD patterns are presented in Fig. 5.4. Black lines in Fig. 5.3 indicate potential profile and the blue lines indicate the current profile. The vertical black lines on the potential profile indicate the points while the in-situ XRD pattern was recorded.

Cycle	Discharge	Charge	Constant	Constant	Rest	Constant	Rest
Number	Rate	Discharge	Current	Voltage	Period	Current	Period
		Potential	Charge	Charge	(min)	Discharge	(min)
1	0.1C	2.5 to 4.1 V	0.3602 mA	4.1 V	15	0.3602 mA	15
2	0.3C	2.5 to 4.1 V	0.3602 mA	4.1 V	15	1.081 mA	15
3	0.5C	2.5 to 4.1 V	0.3602 mA	4.1 V	15	1.801 mA	15
4	1C	2.5 to 4.1 V	0.3602 mA	4.1 V	15	3.602 mA	15

Table 5.1: Protocol of the electrochemical charge-discharge for the in-situ XRD study.



Figure 5.3: Charge discharge cycle at 0.1C rate (1st cycle) during slow rate in-situ Raman spectroscopy. Blue line indicates the current profile and the black line indicate the potential profile and vertical lines on them indicates the point when the Raman spectra were recorded.

Qualitatively, from Fig. 5.4 it is observed that at the beginning of the charging, the peaks of LiFePO₄ were observable on the XRD patterns, while at the end of charging the LiFePO₄ (311) became very weak. At the same time the FePO₄ (311) became dominant indicating the completion of the phase transition of most of the active material particles. It is also found that at the end of the discharge process, LiFePO₄ (311) peak again became strong and FePO₄ (311) weakened.



Figure 5.4: XRD patterns collected during electrochemical charge discharge at 0.1C (1st cycle). R1, and R2 indicates the XRD patterns collected during rest period and C1, C2, etc indicates the XRD patterns taken during charge and D1, D2, etc are XRD patterns taken during discharge as indicated in Fig. 5.4.

To quantify the phase transition process, the peaks in the XRD patterns were deconvoluted using Gaussian fit and the peak intensities were obtained and used to quantify the relative amount of the present phases. The de-convolution of one representative scan of D6 using Gaussian fit peak is shown in Fig. 5.5. Before the deconvolution, the XRD pattern was corrected using a constant baseline. From the deconvoluted peaks, the peak positions, and heights were obtained. Peak positions are almost the same during the course of charging and discharging process. From the resolution of the obtained XRD patterns, it is difficult to draw any conclusion about the lattice strain during Li ion de/intercalation process. As the peak positions were almost identical, it is concluded that during electrochemical charge discharge, only two end members (LiFePO₄ and FePO₄) were present and no compound with intermediate composition (Li_xFePO₄) was observed in the cycles. However, the peak height was used to obtain the relative amount of phases at any stage of electrochemical charge or discharge, which gives the progress of the phase transition.



Figure 5.5: Peak de-convolution and Gaussian fit analysis of one representative XRD pattern D6 collected during electrochemical discharge at 0.1C. The time when D6 is indicated by the 6th vertical line on the potential profile of discharge curve in Fig. 5.3.

Fig. 5.6 a) shows the profile of relative amount of two phases during charge and discharge (Fig. 5.6 b)) of the first cycle. The solid lines with the square symbols (Fig. 5.6 a)) indicate the relative amount of two phases; black lines are for LiFePO₄ and blue lines are for FePO₄. It is observed that at the end of charging, the LiFePO₄ phase was not

completely transformed. In fact, if a longer constant-voltage charging step was used, the phase transition would progress even more toward the complete transformation. But a shorter constant voltage charge period (15 minutes) was actually used. From the Gaussian fit analyses, it was found that at the end of charging during the first cycle, the completed phase transition was approximately 80%. So to get the relative amount of phase transition during discharging, this remaining 20% untransformed materials is deducted from the amount of the second phase (during discharge, the second phase is LiFePO₄). Eventually, the progress profile of the phase transition during 0.1C rate discharge obtained is presented in Fig. 5.7. The dotted line in Fig. 5.7 represents the theoretical path of phase transition progress and it is found that the slow discharge at 0.1C almost follow the equilibrium path. This indicates that during a slow discharge, Li ions get enough time to maintain the equilibrium phase transition while in case of fast discharge the phase transition progress is expected to deviate from the theoretical path.



Figure 5.6: (a) Phase transition progress during electrochemical charge discharge at 0.1C (1st cycle). The dotted lines indicate the equilibrium progress paths. (b)
 Electrochemical charge discharge profiles of the 1st cycle, where blue lines indicate current profile and black lines indicate potential profile.



Figure 5.7: Phase transition progress line during electrochemical discharge at 0.1C (1st cycle) after the correction of history from the charge process. The dotted line indicates the equilibrium progress path.

5.4.2 Effect of Discharge rate

The in-situ XRD patterns obtained during other cycles are presented in Fig. 5.8. From Fig. 5.8 (a) it is found that from scan C4 to R3_2 (the rest period after charge during the 2nd cycle), the noise is significant and the diffraction patterns cannot be quantitatively analyzed. The reason for the noise in those XRD patterns is possible scattering from the electrolyte. During the battery assembly, extra electrolyte of not drained from the cell would scatter the X-ray and cause extra noise. During the 2nd cycle of charging, some of the electrolyte appeared on the top surface which caused scattering of X-ray. This electrolyte scattering problem could not be identified until the rest period after the charge cycle. As soon as the problem was identified, the thin layer of extra electrolyte was removed from the top surface of the electrode and electrolyte scattering disappeared.



Figure 5.8: (a) XRD patterns collected during electrochemical charge discharge at 2nd cycle. (b) at 3rd cycle (c) at 4th cycle. R1, and R2 indicates the XRD patterns collected during rest period and C1, C2, etc indicates the XRD patterns taken during charge and D1, D2, etc are XRD patterns taken during discharge.

XRD patterns during the 2nd, 3rd and 4th cycles of the charging and discharging were analyzed using Gaussian fit analysis and presented in Fig. 5.9, Fig. 5.10 and Fig.

5.11. In Fig. 5.9, during the charging process, some of the data points are missing due to electrolyte scattering as explained before. In fact analyzing charging behavior was never a target of this whole experiment and that is the reason why all the cycles were done at the same charging rate of 0.1C. Even though charging of the battery was conducted at 0.1C rate for all the four cycles, the progress of the phase transition during the charging process looks different as shown in Fig. 5.9, Fig. 5.10 and Fig. 5.11. The reason of this difference can be easily understood if the amount of both phases is considered at the end of each discharge process. For example, at the end of 1st cycle the relative amount of FePO4 is 8.191% indicating that 8.191% of active material was not transformed. Similarly, at the end of 2nd, 3rd and 4th cycle, the untransformed material was 4.41%, 10.79% and 15.78%, respectively. Because these cycles have different pre-history from the previous discharging step, the progress of phase transition during the subsequent charge process looks different in all cycles.

If the battery usage is considered, the battery discharge process is the most important factor as the discharge of the stored energy is what goes to the device. And also although fast charging of the battery is desired, slow charging of the battery still can be tolerated. But considering various applications of the battery systems, some applications require the battery to discharge quickly to the devices and some applications can accept slow discharge of the battery. Fast discharge may lead to the kinetics-induced non-equilibrium phase transition of LiFePO₄. To study this non-equilibrium phase transition, the battery was tested at four different discharge rates such as 0.1C, 0.3C, 0.5C and 1C. However, because of the phase transition history of the previous cycle, it is hard to conceive the real phase transition progress during each discharge (at different discharge rates) from Fig. 5.9, Fig. 5.10 and Fig. 5.11. To get the real quantitative phase transition progress during all the discharge cycles, the amount of untransformed LiFePO₄.

from the charge of the same cycle (taken from the last XRD scan during the rest period after charge) was deducted from the amount of LiFePO₄ during discharge.



Figure 5.9: (a) Phase transition progress line during electrochemical charge discharge during 2nd cycle (0.3C discharge). The dotted lines indicate the equilibrium progress paths. (b) Electrochemical charge discharge profiles of the 2nd cycle, where blue lines indicate current profile and black lines indicate potential profile.

For 1st, 3rd and 4th cycle, the amount of untransformed LiFePO₄ at the end of the rest period after the charge of the respective cycle is 21.19%, 34.85%, and 37.23% respectively. For the 2nd cycle, due to scattering of X-ray by the extra electrolyte, quantitative analysis could not be performed during most of the charge cycle and the rest period after the charge. So, for the 2nd cycle the amount of untransformed LiFePO₄ at the end of the charge and rest period is assumed to be 25%. After correcting the progress of the phase transition curve for all cycles, the progress of phase transition during discharge at different rates are shown in Fig. 5.12.

The dotted diagonal in Fig. 5.12 line indicates the equilibrium phase transition progress during the discharge process. As is already discussed, during slow discharge the phase transition progress is close to the equilibrium path (Fig. 5.7). But the progress of phase transition at higher discharge rates was delayed as the progress line deviated from the equilibrium path. Similar kind of phase transition delays were also predicted by the study in reference [95].



Figure 5.10: (a) Phase transition progress line during electrochemical charge discharge during 3rd cycle (0.5C discharge). The dotted lines indicate the equilibrium progress paths. (b) Electrochemical charge discharge profiles of the 3rd cycle, where blue lines indicate current profile and black lines indicate potential profile.



Figure 5.11: (a) Phase transition progress line during electrochemical charge discharge during 4th cycle (1C discharge). The dotted lines indicate the equilibrium progress paths.
(b) Electrochemical charge discharge profiles of the 4th cycle, where blue lines indicate current profile and black lines indicate potential profile.



Figure 5.12: Phase transition progress line during electrochemical discharge at all four different discharge rates after the correction of history from the charge process. The dotted line indicates the equilibrium progress path.

5.5 In-situ Raman results vs. in-situ XRD results

Phase transition progress was recorded using the in-situ XRD technique for the four different discharge cycles and was also studied at 0.3C discharge rate using the insitu Raman spectroscopy. The results from two techniques were compared in Fig. 5.13.

In chapter 4, the results from the in-situ Raman spectroscopy is discussed and it is concluded that in any particle once the phase transition starts, it finishes very quickly within the same particle. But from the in-situ XRD results (Fig. 5.12) it is observed that the progress of the phase transition is gradual and continuous during the entire discharge process. The reason of this difference of phase transition from two different experiments is that two experiments study the same phenomena at two different length scales.



Figure 5.13: Phase transition progress lines obtained during different discharge rates and from in-situ Raman spectroscopy and from in-situ XRD study.

Raman spectroscopy uses the laser with a spot size of 2 µm which only focuses on a few active materials, and hence can safely be considered as the particle level technique. On the other hand, in-situ XRD focuses on a large area which includes many particles. During the discharge process, among these many particles (X-ray irradiated area) at any point some particles started transforming and the transformation finished very quickly on those particles while the rest of the particles did not start yet. This gives the average result from all the particles (X-ray irradiated area). As the discharge process continued more and more particles started transforming, but once transformation starts in any particle, it finishes very quickly in that particle. As a result, the in-situ XRD study actually investigated the fraction of particles that underwent phase transition, rather than the same process within individual particle. The results show that there was a gradual and continuous phase transformation over the entire period of discharge. This phase transformation mechanism is shown with the help of a simple schematic in Fig. 5.14.



Figure 5.14: Schematic explaining phase transition mechanism during equilibrium and non equilibrium process at slow and high discharge rates. In-situ Raman study area and in-situ XRD study area in the electrodes are also marked in the schematic.

Phase transition delays during fast discharge can also be explained using the same schematic in Fig. 5.14. During fast discharge process, non-equilibrium phase transition was caused by several factors such as, insufficient Li ion availability due to the hindrance of Li ion diffusion through complex porous network, and limited local electronic conductivity. Due to these factors at a high discharge rate, some particles that can start transforming in the slow discharge or in equilibrium could postpone the phase

transformation. The delayed phase transition process is reflected in the in-situ XRD results at higher discharge rates and also in the in-situ Raman results.

5.6 Conclusion

In-situ XRD experiment was designed and employed to study the phase transition mechanism of LiFePO₄ material. Four different discharge rates of 0.1C, 0.3C, 0.5C and 1C were tested and analyzed. The key findings were as follows:

- During charging or discharging processes of all cycles, no noticeable peak shift was observed for both LiFePO₄ (311) and FePO₄ (311) indicating that only two end members (LiFePO₄ and FePO₄) exist and no intermediate composition of Li_{1-x}FePO₄ or Li_xFePO₄ was observed.
- 2) In the XRD experiment a large area of electrode was probed which has many particles, and hence the XRD results indicate the average of phases present in that large X-ray irradiated area. This gives a gradual and continuous progress of phase transition during the entire discharge or charge process.
- Results from the XRD experiments do not contradict with the in-situ Raman results; rather it further validates that in any LiFePO₄ particle once the phase transition starts, it finishes quickly in that particular particle.
- Non-equilibrium phase transition was observed at fast discharge process resulting in a delay in phase transition.
- The in-situ XRD experiments and in-situ Raman experiments are compared and analyzed to establish a phase transition mechanism and pathway.

Chapter 6

Numerical study of phase transition

6.1 Introduction

The phase transformation between the Li-rich phase, β (e.g., LiFePO₄) and the Li-deficient phase, α (e.g., FePO₄) has been demonstrated to play a critical role in the rate performance of cathode materials, such as Mn²⁺-doped LiFePO₄ [55] and nano-size LiFePO₄ [10, 128]. Through in-situ Raman and in-situ XRD experiments, the phase transition phenomena have been already established. The results from those experiments show the presence of phases of nearly stoichiometric end members (LiFePO₄ and FePO₄) during charge/discharge processes and suggest almost no solid solubility between them. The experiments were done using LiFePO₄ of approximately 200 nm and of irregular shape. But reduction of particle size and partial substitution of Fe in LiFePO₄ by Mn²⁺ have been found to shrink the miscibility gap, i.e., to increase the solid solution range. Furthermore, two-phase miscibility gap and ion transport during charge and discharge, depends on temperature [60, 129], composition [130], particle size [10, 130], surface coatings [131], particle shapes [132, 133], and discharge rate [7]. And also fast discharge (insertion) or charge (extraction) may lead to a kinetically-induced non-equilibrium phase transformation which affects the miscibility gap. Recent research also demonstrated that phase transformation in LiFePO₄ is strongly controlled by the strain energy induced from volume difference between the two phases [134]. Additional energy penalty due to the non-equilibrium phase transformation can further decrease the driving energy for phase transformation, resulting in a decrease of the discharge potential but increase of the charge potential [135]. Although these results demonstrate the importance of miscibility gap and non-equilibrium phase transformation, the actual mechanisms and the corresponding implications are not clear.

Phase transition of individual LiFePO₄ particles of different particle size and shape was studied using a numerical study. Here, in this chapter, a unique Mushy-zone approach is introduced and used to study the phase transition behavior of single particle LiFePO₄. Also non-equilibrium phase transition at high discharge rates is studied.

6.2 Mushy Zone Approach

The kinetically-induced non-equilibrium phase transformation can be illustrated by Fig. 6.1. At a low-rate discharge (e.g., C/10 in Fig. 6.1a), a flat plateau appears at around 3.5 V for a LiFePO₄-based cathode, corresponding to the two-phase region in the equilibrium phase diagram (Fig. 6.1b). This equilibrium phase separation, as shown by the red curve in Fig. 6.1b, is realized by Li ion diffusion and readjustments in composition at the two-phase interfaces. However, at a high-rate discharge (e.g., 20 C in Fig. 6.1a) the flat plateau may disappear and is instead replaced with a fast decaying curve [136, 137]. These phenomena are more pronounced at lower temperature where Li diffusion is retarded. It is hypothesized that a non-equilibrium phase transformation occurs during fast discharge, where the non-equilibrium miscibility gap shifts to higher Li contents and is represented by the dashed blue line in Fig. 6.1b. The implication of this phenomenon is that the non-equilibrium miscibility gap is significantly expanded, leading to significant reduction in two-phase solubility and battery discharge capacity.



Figure 6.1: Schematic illustration of equilibrium and non-equilibrium phase transformation in (a) battery discharge and (b) phase diagram. α and β are referring to the Li-poor phase (FePO₄) and Li-rich phase (LiFePO₄), respectively.

Though the well-known "core-shrinking" description [75, 138] has been widely employed to simulate the phase separation between the two solid solution regions outside the miscibility gap; however, it fails to account for the non-equilibrium phenomenon. An approach termed as "mushy-zone" (MZ) approach, has been developed as an effective tool to study the origin and implication of the non-equilibrium phenomenon and also the miscibility gap dynamics during charge/discharge process and particle size and morphology effect.

6.2.1 Model Description

Accounting for non-equilibrium effect.— The rate of phase transformation depends on the interface mobility and Li diffusivity in both the α - and β -phases. Studies conducted by Wang et al.[139, 140] have confirmed the importance of the mixed mode of phase transformation, where the interface mobility is much slower than the Li-ion diffusion in the individual phase. Experimentally, high-resolution TEM images showed that Li insertion and extraction seem to be concentrated at the phase boundary, with

negligible transfer occurring in either the α - or β -phase [42]. Therefore, one possible approach to account for the effect of non-equilibrium phase transformation is to simply implement a sluggish interfacial diffusion coefficient in the continuum diffusion equation. The derivation of the interfacial diffusion coefficient is detailed in the following paragraphs.

Simulation of battery discharge.— Similar to the approach employed by Wang et al, [139, 140] only 1D Li transport within a single particle (plate or sphere) is simulated. The governing equations are

$$\frac{\partial c_{Li}}{\partial t} = D_{Li} \frac{\partial^2 c_{Li}}{\partial x^2}$$
 for plate shaped particle......6.1

$$\frac{\partial c_{Li}}{\partial t} = \frac{1}{r^2} \frac{\partial}{\partial r} \left(D_{Li} r^2 \frac{\partial c_{Li}}{\partial r} \right) \quad \text{for spherical particle......6.2}$$

with the following boundary conditions

Particle center x = 0 or r = 0 $j_{Li} = 0$

Particle surface
$$x = l_0$$
 or $r = r_0$ $j_{Li} = \frac{l_n}{F}$ 6.3

.

where j_{Li} s the Li flux and i_n is the intercalation current.

Ignoring the anode overpotential and electrolyte resistance, the discharge potential E is determined from the discharge kinetics, i.e.,

where *k* is a lumped factor accounting for both the exchange current and electrolyte concentration, c_s Li concentration at the particle surface, and *F* Faraday constant.

The implementation of the diffusion coefficient within the miscibility gap (either equilibrium or non-equilibrium) is straightforward. The overall Li flux j_{Li} in one-dimension within the two-phase region can be described by

where b is the volume fraction of the β -phase, which is determined by the lever rule.

Within the miscibility gap, the overall Li concentration $\,c_{{\scriptscriptstyle L}i}\,$ is

If at thermodynamic equilibrium, compositions of the two phases at the miscibility gap have the following relationship

Combining Eqns. 6.7 and 6.8, the phase compositions at the miscibility gap can be described by

$$c_{Li,\alpha} = \frac{c_{Li}}{1 - b + bK_{Li}}, \qquad c_{Li,\beta} = \frac{K_{Li}c_{Li}}{1 - b + bK_{Li}}.....6.9$$

The Li diffusion flux in the miscibility gap is thus rearranged by substituting Eqn. 6.9 into Eqn. 6.5 as

Rearranging Eqn. 6.10 and comparing with the Fick's first law result in the following twophase diffusion coefficient of Li

where $f(c_{Li})$ is a tunable function accounting for slow interfacial diffusion.

Parameters used in the simulations are listed in Table 6.1. A commercial computational fluid dynamics (CFD) software was used to simulate the single particle phase transition for different particle size and morphology. The codes used in the CFD software to describe the equations and properties of the material are described in Appendices.

Parameter	Value
Li diffusivity in β phase, $D_{Li,\alpha}$ (m ² /s)	8×10^{-18}
Li diffusivity in α phase, $D_{Li,\beta}$ (m ² /s)	4.8×10^{-17}
Equilibrium Li composition in α phase, $c^{0}_{Li,\alpha}$ (mol/m ³)	616.5
Surface concentration of Li composition at β phase, $c^{0}_{Li,\beta}$ (mol/m ³)	19409.4
Partition coefficient of phase β to phase α , K_{Li}	31.48
Maximum Li concentration, c_{max} (mol/m ³)	22834.6
Cathode kinetic constant, k	3.14×10 ⁻⁶
Tunable function $f(c_{Li})$ in Eq. 6.11	$Ab^{1.8}$
Equilibrium potential of LiFePO ₄ , E_0 (V)	$3.4245 + 0.85e^{-800\left(\frac{c_{s}}{c_{max}}\right)^{-1.3}} - 17e^{-0.98\left(\frac{c_{s}}{c_{max}}\right)^{-14}}$

Table 6.1: Parameters used in the simulation (Refs. [42, 105])

6.2.2 Effect of Diffusivity through two-phase interface

The impact of interfacial Li diffusivity on discharge performance of a sphere is further explored in Fig. 6.2 and Fig. 6.3. At discharge rate of 2C, low Li diffusivity makes the transition between the stages 2 and 3 less pronounced as shown in the inset of Fig. 6.2, indicating significant delay in phase transition. With the increase of Li diffusivity from $A \sim 10^{-4}$ (*A* is a tunable factor), the discharge capacity first jumps from $\sim 80 \text{ mAhg}^{-1}$ to more than 120 mAhg⁻¹, and then gradually increases and eventually levels off at around 160 mAhg⁻¹ when *A* reaches $\sim 10^{-3}$. At discharge rate of 20C in Fig. 6.3, the slope of the potential flat plateau as seen during low-rate discharge is declining as demonstrated from both the discharge performance curves and the surface composition profiles. When the discharge reaches the stage 2, the surface composition continuously increases till the end of discharge. This is more pronounced for slow interfacial diffusion.



Figure 6.2: Impact of interfacial diffusion coefficient on discharge capacity and battery performance (inset) for a sphere with radius of 20 nm. Values of *A* in the function $f(c_{II}) = Ab^{1.8} \text{ are } \mathbf{1}: 9.45 \times 10^{-5}, \mathbf{2}: 1.08 \times 10^{-4}, \mathbf{3}: 1.35 \times 10^{-4}, \mathbf{4}: 6.74 \times 10^{-4}, \text{ and } \mathbf{5}: 1.08 \times 10^{-3}.$



Figure 6.3: Impact of interfacial diffusion coefficient on discharge performance and particle surface composition (inset). The simulation is done for a sphere with radius of 20 nm at discharge rate of 20C. The dash lines in the inset indicate the equilibrium and

maximum compositions.

6.2.3 Validation of model using equilibrium simulation

To validate the suitability of the MZ approach the equilibrium phase transformation is studied under low discharge rate and compared with core-shrinking model. This would be a baseline for further study of non-equilibrium phase transformation. The 1D simulation result of a plate at low discharge rate (0.5C) which is believed to be equilibrium phase transformation is shown in Figure 6.4 which compares the Li distribution simulated by the MZ approach with those predicted by the coreshrinking formulation (adopted from Zhang and White [38]). The sharp interfaces and moving boundaries have been successfully predicted by the MZ approach. Figure 6.4

indicates that the two approaches essentially predict almost identical phase compositions and interface locations. Also variation of Li ion distribution at the front and back surface of the plate obtained from two different approaches is compared in Figure 6.5. Here also results from two approaches show very good agreement. Three distinct stages labeled as 1, 2, and 3 in the Figure 6.5 are in agreement with the well-known three stages of phase separation during the discharge. The stage '2' corresponds to the two-phase transformation, and the stages '1' and '3' are for the single-phase transformation either in the α - or β -phases.



Figure 6.4: Li concentration distribution along the normalized thickness of a 20 nm-thick plate at discharge rate of 0.5C. *A*=0.00135 in $f(c_{Li})$ has been used for the MZ model.



Figure 6.5: Variation of Li concentration at the front and back surfaces of the 20 nm-thick plate during discharge at 0.5C. The inset compares the discharge performance.

6.2.4 Particle Geometry Effect

As the suitability of the MZ approach has been established, the next step is to study the effect of particle shape under non-equilibrium phase transformation. Particle shape is important in phase transformation, since it affects the directions of Li diffusion and moving boundaries. For example, the phase boundary moves unidirectionally in a plate, while convergently toward the center in a sphere. The impact of particle shape is demonstrated in Fig. 6.6. Particles with different shapes show similar discharge performance up to 120 mAhg⁻¹, beyond which the sphere experiences a dramatic drop in cell voltage while the plate extends the discharge capacity to ~162 mAhg⁻¹. A sharp transition from the stage 2 to stage 3, identified from the surface concentration curve, is

observed for the plate but absent for the sphere. This can be explained by different Li flux across the two-phase interface for different particle shapes, i.e.,

Sphere:
$$J = \frac{L^3 \rho Q_{th} C_{rate}}{3r^2}$$
 Plate: $J = L \rho Q_{th} C_{rate}$ 6.12

where Q_{th} is the theoretical capacity, r is the radius of the α -phase core, and L is the radius of the sphere or thickness of the plate ($L = r_0 = l_0$). For the spherical particle, at the beginning of discharge (r = L), the Li flux at the particle surface is only one third of that of the plate. However, as discharge proceeds, the two-phase interface between the α - and β -phases shrinks dramatically in the sphere. As a result, the Li flux increases rapidly as a function of r^2 (Eqn. 6.12). At r = 0.58L, the two-phase interfacial flux in the sphere equals that in the plate. Upon further discharging, the Li flux across the two-phase boundary may not keep up with the required discharge rate owing to the sluggish Li diffusivity, thus developing the non-equilibrium phenomenon. This generates the gradual increase of the surface composition for the spherical particle in Fig. 6.6. On the other hand, the two-phase interfacial flux in the plate-shaped particles remains constant throughout the entire discharge process.

For higher discharge rate at 2C, Fig. 6.7 clearly indicates that delicate non-equilibrium phase transformation can be brought about by the kinetically-induced process. At the relatively higher discharge rate, Li ion diffusion and readjustments at the miscibility gap may not keep up with the high flux of Li ions from the particle surface. Therefore, equilibrium phase transformation is destroyed and non-equilibrium behavior appears. For a 20 nm-thick plate in Fig. 6.7a, the miscibility gap is found to expand, especially toward the β -phase. The relative expansion ratio of the two-phase miscibility gap is calculated to be 0.8% and 6% at 1C and 2C discharge rates, respectively. Though the expansion ratio at 2C discharge rate is large; however, it might seem that the simulation on a 1D plate

exaggerates the effect of non-equilibrium phase separation, since the Li inflow rate at a plate surface is 3 times of that at a sphere surface (when $r_0=I_0$) at the same discharge rate.



Figure 6.6: Discharge performance of a $LiFePO_4$ plate and a sphere at 0.5C. The inset shows particle surface concentration of Li near completion of discharge.

For a spherical particle in Fig. 6.7b, due to variation of the two-phase interfacial flux during discharge, the sphere first experiences equilibrium phase transformation at early stage of discharge and then non-equilibrium phase transformation upon further discharging. The β -phase composition at the miscibility gap stays at the equilibrium value initially and then increases gradually with the state of discharge (SOD). SOD is defined as the ratio of discharging capacity at a specific point to the full capacity. Near the end of discharge (e.g., SOD=0.833 when r<0.58 L), the miscibility gap expands promptly when significant non-equilibrium phenomenon develops.



Figure 6.7: Li concentration profiles along the normalized thickness during discharge in
 (a) a plate (thickness l₀=20 nm) and (b) a sphere (radius r₀=20 nm) at 2C discharge rate.
 The state-of-discharge is indicated in the figures. The dash lines indicate the thermodynamic equilibrium phase compositions.

The different rate of phase transformation could lead to different discharge behavior of these particles. Interestingly, Fig. 6.8 shows that the curves of capacity as a function of discharge rate for plates and spheres cross at 1.5C. At lower rates, plateshaped particles are slightly better but show much lower capacity when discharge rate increases.



Figure 6.8: Comparison of capacity as a function of discharge rate for plate-shaped and spherical particles.

The interesting observation in Fig. 6.8 can be interpreted by different Li flux across the two-phase interface for different particle shapes [7]. The interfacial Li flux for plate shaped particles remains constant at a discharge rate throughout the entire discharge process, while the flux increases for spherical particles as phase transformation approaches the center of the particles. The ratio of interfacial flux between the two shapes shown in Fig. 6.9 indicates that at the early stage of phase transformation (r/L > 0.58 where r is the radius of the FePO₄ core and L is the plate thickness), the Li flux (required by the discharge rate) across the two-phase interface is notably smaller for spherical particles; however, as the discharge proceeds till r/L < 0.58 the interfacial Li flux in spheres drastically increases due to significant shrinkage of the two-phase interface, and eventually causes the kinetically-induced non-equilibrium phase transformation. As

the interfacial flux spikes near the end of discharge leading to non-equilibrium phase transformation, the core of a LiFePO₄ sphere could not be utilized even at an extremely small rate. However, a LiFePO₄ plate could be completely accessed as long as the discharge rate is low enough, resulting in a higher discharge capacity. The ratio of r/L = 0.58 translates to a discharge capacity of $170 \times (1^3 - 0.58^3) \approx 140$ mAh/g for a LiFePO₄ sphere, where 170 is the theoretical capacity (in mAh/g) of LiFePO₄ materials. When the discharge rate increases, capacity of LiFePO₄ particles in both shapes decreases. However, when the discharge capacity decreases to below 140 mAh/g (r/L = 0.58), the lower interfacial Li flux as seen in Fig. 6.9 makes the sphere better utilized and therefore results in a higher discharge rate. The capacity of 140 mAh/g (r/L = 0.58) should therefore correspond to the point where the two curves for plate and sphere cross as shown in Fig. 6.8.



Figure 6.9: Ratio of the Li ion flux at the two-phase interface for two different particle shapes during discharge.
6.2.5 Particle Size Effect

Apart from the significant effect of particle geometry, the effect of particle size also has been studied. Discharge performance of LiFePO₄ spherical particles with different sizes was numerically simulated and is shown in Figure 6.10. At low rate (0.5C), capacities of spherical particles with different sizes are almost the same, but at higher rates (1C and 2C) bigger particles show significantly lower capacity. For example, the 20 nm LiFePO₄ particle delivers a capacity of 160 mAh/g almost constantly at different discharge rates; however, the 50 nm particle only shows a capacity of 100 mAh/g at a discharge rate of 2C though its capacity is around 160 mAh/g at smaller rate of 0.5C. The fluctuation in some of the discharge curves at high discharge rates or for bigger particles is due to the fact that the potential plateau is calculated based on single particle. The discharge curves we get experimentally are usually from a many particle electrode system where the potential value is an average value from many particles. But for a single particle, at higher discharge rates or bigger particles use the fluctuation of Li ions, and available Li ions in the adjacent electrolyte cause the fluctuations in the potential plateau.

The influence of particle size on discharge curve at different rates can be explained by the Li flux at the two-phase interface between α (lithium-deficient phase, FePO₄) and β (lithium-rich phase, LiFePO₄) phases. In Eqn 6.12, r/R (L=R) can be considered as the state of discharge (SOD); a smaller value of r/R ratio translates to a larger SOD (close to the end of discharge). It is evident from Eqn 6.12 that for bigger particles, the lithium flux at the two-phase interface is higher at the same SOD and it becomes more prominent near the end of discharge.



Figure 6.10: Discharge performances of single spherical particle of $LiFePO_4$ at a) 0.5C, b) 1C and c) 2C rate.

Size of the particles also impacts the miscibility gap of the two phase region in the phase diagram. This size dependence of the miscibility gap during the discharge process is shown in Figure 6.11. For a fixed particle size at a constant discharge rate, the miscibility gap slowly increases during the discharge process and increases more rapidly near the end of the discharge. Also, it is evident that the miscibility gap expands more significantly for bigger particles. At lower discharge rate, the miscibility gap does not expand significantly with the particle size. But at higher discharge rate, a little increase in the particle size causes a large expansion in the miscibility gap of the two-phase region. This implies that for bigger particles of active materials, non-equilibrium phase transition can occur at even a low discharge rates. This clearly indicates that non-equilibrium miscibility gap not only relies on temperature, discharging or charging time etc., but also on particle size. This is due to the fact that near the end of the discharge process the sluggish lithium diffusion cannot keep up with the increasing lithium flux at the two-phase interface demanded by the prescribed discharge rate. As a result, the equilibrium lithium concentrations in the two phases across the boundary cannot be maintained; therefore, the kinetics-controlled non-equilibrium phase transformation occurs and the miscibility gap expands.



Figure 6.11: Miscibility gap during the discharge process for different particle size at a)

0.5C, b) 1C and c) 2C discharge rate.

6.2.6 Nature of two-phase interface

The nature of the two-phase interface can be understood from the concentration profile at any SOD. At a SOD of 44%, the concentration profiles are shown in Figure 6.12 for two different discharge rates. For lower discharge rates, the interface is wider than that at higher discharge rates, especially for smaller particle size, e.g., 20 nm at 0.5C. And also at the same discharge rate, bigger particles show much narrower two-phase transition. The results indicate that small particle size (<20 nm) and slow discharge rate (<0.5C) tend to create homogenous phase transformation (i.e., small miscibility gap). The results of this study are in good agreement with Meethong et al's [105] experimental discovery that miscibility gap contracts systematically (increased solid solution limits) with decreasing particle size. The results in Figure 6.12 can also be explained from material structure point of view. During the discharge process, when lithium ions enter the active material particles, there comes a mismatch between the de-lithiated α (FePO₄) and lithiated β (LiFePO₄) phase. When the particle size is smaller, this mismatch is accommodated within several layers of FeO₄ blocks [73]. But when the particle size is bigger, the mismatch between the two phases is accommodated by creation of dislocations and cracks. So, for smaller particles like particles of 20 nm radii show wider two-phase interface, while the 40 nm and 50 nm particles show very sharp two-phase interface.



Figure 6.12: Li concentration profiles at SOD of 44% for different particle size at a) 0.5C,

b) 1.0C, and c) 2.0C.

To understand the nature of sharp two-phase interface at any stage of discharge,

the gradient of chemical potential $\begin{pmatrix} d\mu \\ dr \end{pmatrix}$ for Li ion intercalation is calculated for

different particle sizes. The following Cahn-Hilliard equation was used for 1D Li ion transport [141].

$$\nabla \mu_{Li} = \left\{ \left[-2W_c + RT \left[\frac{1}{c} + \frac{1}{(1-c)} \right] \right] / V_m + C_{1111} \cdot \Delta e_{11}^0 \cdot \Delta e_{11}^0 + C_{2222} \cdot \Delta e_{22}^0 \cdot \Delta e_{22}^0 + C_{3333} \cdot \Delta e_{33}^0 \right\} \cdot \frac{\partial c}{\partial x} + \left(2C_{1133} \cdot \Delta e_{33}^0 - 2C_{1122} \cdot \Delta e_{11}^0 \cdot \Delta e_{22}^0 - 2C_{2233} \cdot \Delta e_{33}^0 \right) \frac{\partial c}{\partial x} - \kappa \cdot \left(\frac{\partial^3 c}{\partial x^3} \right) \dots 6.13$$

where V_m is the molar volume of stoichiometric LiFePO₄, W_c is the regular solution coefficient, C_{ijkl} is the stiffness tensor component, Δe_{ij}^{o} is linear misfit between the two end members (LiFePO₄ and FePO₄), and κ is the concentration gradient coefficient. The gradient of chemical potentials for three particle sizes is plotted in Fig. 6.13.



Figure 6.13: Gradient of chemical potentials of spherical particles at 1C discharge rate at 44% SOD (corresponding to Fig. 6.12b).

It is clearly observed that for bigger particle sizes, the change in chemical potential is more drastic than those of smaller ones which is believed to be caused by the sharp two phase interfaces as one sees in Figure 6.12. From a material point of view, the simulated Li concentration in Figure 6.12 and chemical potential gradient shown in Figure 6.13 are determined by the mismatch of structural properties (e.g., structure and lattice parameters) of the two phases presented in the cathode materials.

6.3 Conclusion

The important non-equilibrium phenomena that have pronounced effects on battery discharge performance have been explored, via the unique mushy-zone approach. The non-equilibrium phase transition, occurring due to the fact that sluggish Li diffusion across the two-phase boundaries cannot keep up with the required Li flux at high discharge rate. As a result, the two-phase region in the phase diagram expands and the miscibility gap shifts towards the Li-rich phase. Additionally, particle shape is found to significantly affect the non-equilibrium phase transformation due to different Li flux across the two-phase boundary and therefore different degree of non-equilibrium phase transformation. Growth of LiFePO₄ phase follows a well-predicted linear curve in plate-shaped particles; however, two-phase transformation is remarkably delayed for spherical particles, especially at low discharge rates. The different rate: at lower rates, plate-shaped particles are slightly better but show much lower discharge capacity when discharge rate increases; at higher discharge rates, spherical particles show better capacity.

Non-equilibrium phase transformation tends to be pronounced for bigger particles even at low discharge rates. It is found that the miscibility gap expands gradually during discharge, which is evident for bigger particles, especially at high rates. Internal distribution of Li within a particle suggests that small particles (20 nm) tend to create homogenous phase transformation (i.e., no miscibility gap).

It is admitted that experimental validation of the simulation results on individual particle in this work is almost impossible. The conventional battery testing picks up the electrode-level phase transformation information which is thus averaged over multiple particles. Besides, in a complex multi-phase electrode structure the lithiated and delithiated phases may coexist within the same particle or as an assembly of particles [132]. Nevertheless, this study is intended to understand the non-equilibrium phase transformation in individual particle of LiFePO₄ and provide fundamental guidance to design high-rate battery materials from a particle-level point of view.

Chapter 7

Conclusion

LIBs have become very popular over the last two decades because of high energy densities in various applications including electric vehicle and portable electronic devices. LiFePO₄ material has become one of the most reliable and high performing cathode materials due to its inherent stability and safety features. The phase transition between the two end members (Li rich LiFePO₄ and Li poor FePO₄) plays a critical role in the performance of the LIB containing this material. And hence in this work this material has been studied to understand their two-phase phase transition behavior.

LiFePO₄ material has been synthesized using solid state synthesis method. Material synthesized at 700 °C showed an irregular morphology with an average particle size of approximately 200 nm. Phase has been confirmed to be olivine type structure (*Pnma*). Two unique techniques, i.e., in-situ Raman spectroscopy and in-situ XRD, were designed to study the phase transition phenomena in LiFePO₄. In-situ Raman studies target the particle-level phenomena while the in-situ XRD studies focus on the phase transformation of the materials at the bulk or electrode level.

A specially designed transparent cell containing a NaCl window was used for the in-situ Raman study. Raman laser spot size was 2 µm while our particle size was approximately 200 nm, which indicates that the Raman laser focused only on a few particles. The particle-level Raman study of LiFePO₄ at a slow discharge rate (0.3C) showed that during both charge and discharge the particles under study showed a quick phase transition which is agreement with the domino-cascade model [73]. At the particle level, once the phase transition starts in any LiFePO₄ particle, it finishes very quickly while some other particles may not have started the phase transition yet. Thus at any

stage of charge or discharge of a LIB, Raman laser detects either transformed particles or the particles those have not been transformed. At high discharge rates, a phase transition delay has been observed. At high discharge rates, the same particles that show phase transition at slow discharge rates may be completely suppressed and may not go through phase transition depending on the local electronic conductivity and availability of Li ions through porous electrode.

Similarly, the other specially-designed transparent cell that was used in the insitu XRD equipment was developed. A uniquely-designed electrode containing LiFePO4 was used to make batteries for the in-situ XRD experiments. The battery was tested for four different discharge rates of 0.1C, 0.3C, 0.5C and 1C. For these discharge rates, the phase transition progress curves showed a gradual increase during the entire period of cell discharge. This apparently contradicts the in-situ Raman results. But the in-situ XRD results were discovered on a larger area of a battery electrode that contains lots of active material particles, while the Raman laser probes a very small area that contains only a few particles. Although the electrode structure is heterogeneous, and electronic conductivity, electrolyte porous channels, and other electrochemical parameters can vary throughout the electrode, the electrode is quite homogenous over a small area like 2 µm (the size of the Raman laser spot). Therefore, the particles under the laser during the insitu Raman study showed either transformed or non-transformed. But for a larger area (XRD study area) that was composed of transformed and non-transformed particles, with the course of discharge the amount of transformed particles increases which gives a gradually and continually increasing phase transition progress curve So, this gradual increase of phase transition is in agreement with the core-shrinking model [75] and indicates that core-shrinking model only works at the bulk level, or the aggregate level. In-situ XRD study also revealed that there was almost no detectable peak position shift for LiFePO₄ and FePO₄ phases indicating the presence of only stoichiometric end members (LiFePO₄ and FePO₄) and no solid solution with the Li_xFePO_4 composition. Besides, the in-situ XRD results also showed that the phase transition was delayed at high discharge rates due to the kinetics-induced non-equilibrium phase transition.

Numerical study using the Mushy Zone approach on individual LiFePO₄ particles was performed to investigate the effect of particle size and morphology on phase transition, as well as the miscibility gap dynamics during equilibrium and non-equilibrium phase transition at slow and high discharge rates. It is found that at low discharge rates, the plate shaped particles show higher capacity than spherical particles and at high discharge rate beyond a critical discharge rate, the spherical particles show higher capacity. For spherical particles, the miscibility gap increases exponentially near the end of the discharge process. Particle size of spherical particles plays a key role in phase transition. With the increase of particle size, the miscibility gap becomes wider which makes difficult for the phase transition to be completed. Two phase interface for smaller particles is wide while the large particles show a very sharp and narrow two-phase interface. Large particle can show non-equilibrium behavior at the same discharge rate that gives the equilibrium situation in a small particle.

This study provides a detailed picture of phase transition phenomena in LiFePO₄ through both the in-situ Raman and XRD studies that record the phase transition information in real time. And also the numerical study opens up the opportunity to study the same phenomena in a single particle. The results obtained from studies are expected to help to understand the complicated phase transition behavior of LiFePO₄ which will eventually shed lights on building high performing and reliable LIBs.

Appendix A

Code for simulating Spherical particles

include "udf.h"
include "id.h"

include "mem.h"

include "sg.h"

#define da 8.0e-18 /* diffusivity in alpha phase, LiFePO4 phase */
#define db 4.8e-17 /* diffusivity in beta phase, FePO4 phase */
#define xb0 616.5352 /* surface concentration at beta phase */
#define xa0 19409.44 /* surface concentration at alpha phase */
#define K 31.48148148 /* partition coefficient of phase 2 to phase 1*/
#define I 20.0e-9 /* particle radius*/
#define liflux 0.5 /* unit C */
#define cmax_c 22834.638 /* max surface concentration cathode calculated assuming
170mAh/g */
#define kb 0.00000314 /* Cathode kinetic constant */
#define R 8.314 /* Universal gas constant */

#define T 298.15 /* Room temperature in absolute scale */

#define F 96485.0 /* Faraday constant */

/* C_UDSI(c,t,0) Li total molar fraction */

DEFINE_DIFFUSIVITY(Li_diff,c,t,i)

{

double diff,tt;

double x[ND_ND];

C_CENTROID(x,c,t);

```
/* determining volume fraction of phase 2, with smaller D */
if (C_UDSI(c,t,0)<xb0) diff=db*x[0]*x[0];
else if (C_UDSI(c,t,0)>xa0) diff=da*x[0]*x[0];
else
{tt=(C_UDSI(c,t,0)-xb0)/(xa0-xb0);
diff=((1.0-tt)*db+K*da*tt)*(1.0/(1.0-tt+K*tt))*x[0]*x[0]*0.00235*pow(tt,1.8);
}
return diff;
}
DEFINE_PROPERTY(cell_density,c,t)
{
double rho;
double x[ND_ND];
C_CENTROID(x,c,t);
rho=x[0]*x[0]*l*l;
return rho;
}
/************* calculate discharge curves ****************/
DEFINE_PROFILE(left_in_flux, t, i)
{
cell_t c;
 face_t f;
       begin_f_loop(f,t)
       {F_PROFILE(f,t,i) = 1.0/3.0*liflux/1000.0*170.0*3.6*1000000.0/96485.0*l*l;}
```

```
end_f_loop(f,t)
```

```
DEFINE_EXECUTE_AT_END(cal_profiles)
```

{

double x[ND_ND],xx;

Domain *d;

Thread *t;

cell_t c;

FILE *f1;

int i;

real sum=0.;

f1=fopen("profile.txt","a");

```
fprintf(f1,"%12.8f\n", CURRENT_TIME);
```

```
d = Get_Domain(1);
```

thread_loop_c(t,d)

```
{
```

```
begin_c_loop(c,t)
```

{C_CENTROID(x,c,t);

```
if (c==0) {xx=x[0]*0.0;}
```

fprintf(f1,"%12.8e %12.8e\n",(x[0]-xx), C_UDSI(c,t,0));}

end_c_loop(c,t)

```
fprintf(f1,"\n");
```

```
}
```

```
fclose (f1);
```

}

DEFINE_EXECUTE_AT_END(cal_dis_curves)

{

double yy,Ueq,z,u,y;

Domain *d= Get_Domain(1); /* look up through Boundary panel*/

Thread *t_left = Lookup_Thread(d,5);

int f=0;

FILE *f1;

f1=fopen("dis_curve.txt","a");

```
yy=F_UDSI(f,t_left,0)/cmax_c;
```

```
Ueq=3.4323-5.8428*exp(-80.2493*pow((1-yy),1.3198))-3.2474/1000000.0*exp(20.2645*pow((1-
```

```
yy),3.8003))+3.2482/1000000.0*exp(20.2646*pow((1-yy),3.7995));
```

```
y=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l/(2.0/F*kb*pow((cmax_c
```

```
F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5));
```

 $z = log(y + pow(y^*y + 1.0, 0.5));$

u=z*R*T/0.5/F+Ueq;

fprintf(f1,"%12.8e %12.8e\n",CURRENT_TIME, u);

fclose (f1);

}

```
/******************************* calculate surface concentration curves **********************/
```

```
DEFINE_EXECUTE_AT_END(surf_con)
```

{

```
Domain *d= Get_Domain(1); /* look up through Boundary panel*/
```

```
Thread *t_left = Lookup_Thread(d,5);
```

face_t f;

FILE *f1;

```
f1=fopen("surf_con.txt","a");
begin_f_loop(f, t_left)
{fprintf(f1,"%12.8e
                       %12.8e\n",CURRENT_TIME, F_UDSI(f,t_left,0));}
end_f_loop(f,t_left)
fclose (f1);
}
/************************* calculate surface concentration curves ****************/
DEFINE_EXECUTE_AT_END(core_con)
{
Domain *d= Get_Domain(1); /* look up through Boundary panel*/
Thread *t_left = Lookup_Thread(d,3);
face_t f;
FILE *f1;
f1=fopen("core_con.txt","a");
begin_f_loop(f, t_left)
{fprintf(f1,"%12.8e
                       %12.8e\n",CURRENT_TIME, F_UDSI(f,t_left,0));}
end_f_loop(f,t_left)
fclose (f1);
```

```
}
```

Appendix B

Code for simulating Plate Shaped particles

include "udf.h"

include "id.h"

include "mem.h"

include "sg.h"

#define da 8.0e-18 /* diffusivity in alpha phase */

#define db 4.8e-17 /* diffusivity in beta phase */

#define xb0 616.5352 /* surface concentration at beta phase */

#define xa0 19409.44 /* surface concentration at alpha phase */

#define K 31.48148148 /* partition coefficient of phase 2 to phase 1*/

#define | 20.0e-9 /* particle length*/

#define liflux 0.5 /* unit C */

#define cmax_c 22834.638 /* max surface concentration cathode calculated assuming

170mAh/g */

#define kb 0.00000314 /* Cathode kinetic constant */

#define R 8.314 /* Universal gas constant */

#define T 298.15 /* Room temperature in absolute scale */

#define F 96485.0 /* Faraday constant */

```
/* C_UDSI(c,t,0) Li total molar fraction */
```

DEFINE_DIFFUSIVITY(Li_diff,c,t,i)

{

double diff,tt;

/* determining volume fraction of phase 2, with smaller D */

```
if (C_UDSI(c,t,0)<xb0) diff=db/l/l;
```

```
else if (C_UDSI(c,t,0)>xa0) diff=da/l/l;
```

else

```
{tt=(C_UDSI(c,t,0)-xb0)/(xa0-xb0);
```

```
diff=((1.0-tt)*db+K*da*tt)*(1.0/(1.0-tt+K*tt))/I/I*0.00235*pow(tt,1.8);
```

}

```
return diff;
```

}

```
/****** calculate discharge curves ****************/
```

```
DEFINE_PROFILE(left_in_flux, t, i)
```

```
{
```

```
cell_t c;
```

face_t f;

```
begin_f_loop(f,t)
```

```
{F_PROFILE(f,t,i) = liflux/1000.0*170.0*3.6*1000000.0/96485.0;}
```

```
end_f_loop(f,t)
```

```
}
```

```
DEFINE_EXECUTE_AT_END(cal_profiles)
```

{

```
double x[ND_ND],xx;
```

Domain *d;

Thread *t;

cell_t c;

FILE *f1;

int i;

```
real sum=0.;
f1=fopen("profile.txt","a");
fprintf(f1,"%12.8f\n", CURRENT_TIME);
d = Get_Domain(1);
thread_loop_c(t,d)
{
begin_c_loop(c,t)
{C_CENTROID(x,c,t);
if (c==0) {xx=x[0];}
fprintf(f1,"%12.8e %12.8e\n",(x[0]-xx), C_UDSI(c,t,0));}
end_c_loop(c,t)
fprintf(f1,"\n");
}
fclose (f1);
}
/******************************* calculate discharge curves ******************/
DEFINE_EXECUTE_AT_END(cal_dis_curves)
{
real yy,Ueq;
Domain *d= Get_Domain(1); /* look up through Boundary panel*/
Thread *t_left = Lookup_Thread(d,3);
face_t f;
int n,js,m,i;
double z,y,z1,y1,z0,y0,xn;
double a,b,h,eps;
```

FILE *f1; a=2.0; b=5.0; h=0.1; eps=0.0000001; m=1; f1=fopen("dis_curve.txt","a"); begin_f_loop(f, t_left) yy=F_UDSI(f,t_left,0)/cmax_c; Ueq=3.4323-5.8428*exp(-80.2493*pow((1-yy),1.3198))-3.2474/1000000.0*exp(20.2645*pow((1yy),3.8003))+3.2482/1000000.0*exp(20.2646*pow((1-yy),3.7995));

n=0; z=a;

{

y=liflux/1000.0*170.0*3.6*1000000.0/96485.0*I-1.0/F*kb*pow((cmax_c-

```
F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z-Ueq))-exp(-
```

0.5*F/R/T*(z-Ueq)));

```
while ((z<=b+h/2.0)&&(n!=m))
```

{ if (fabs(y)<eps)

{ n=n+1; xn=z;

z=z+h/2.0; y=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l-

1.0/F*kb*pow((cmax_c-

F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z-Ueq))-exp(-0.5*F/R/T*(z-Ueq)));

}

else

{ z1=z+h; y1=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l-1.0/F*kb*pow((cmax_c-F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z1-Ueq))-exp(-

0.5*F/R/T*(z1-Ueq)));

if (fabs(y1)<eps)

{ n=n+1; xn=z1;

z=z1+h/2.0; y=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l-

1.0/F*kb*pow((cmax_c-

F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z-Ueq))-exp(-

0.5*F/R/T*(z-Ueq)));

```
}
else if (y*y1>0.0)
{ y=y1; z=z1;}
```

else

{ js=0;

while (js==0)

{ if (fabs(z1-z)<eps)

{ n=n+1; xn=(z1+z)/2.0;

z=z1+h/2.0; y=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l-

1.0/F*kb*pow((cmax_c-

F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z-Ueq))-exp(-

0.5*F/R/T*(z-Ueq)));

```
js=1;
}
else
```

{ z0=(z1+z)/2.0; y0=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l-

1.0/F*kb*pow((cmax_c-

F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z0-Ueq))-exp(-

0.5*F/R/T*(z0-Ueq)));

if (fabs(y0)<eps)

{ xn=z0; n=n+1; js=1;

z=z0+h/2.0; y=liflux/1000.0*170.0*3.6*1000000.0/96485.0*l-

1.0/F*kb*pow((cmax_c-

F_UDSI(f,t_left,0)),0.5)*pow(F_UDSI(f,t_left,0),0.5)*(exp(0.5*F/R/T*(z-Ueq))-exp(-0.5*F/R/T*(z-Ueq)));

```
}
          else if ((y*y0)<0.0)
           { z1=z0; y1=y0;}
          else { z=z0; y=y0;}
         }
       }
      }
    }
  }
fprintf(f1,"%12.8e
                %12.8e\n",CURRENT_TIME, xn);
} end_f_loop(f,t_left)
fclose (f1);
}
DEFINE_EXECUTE_AT_END(cal_interface)
```

```
double min, max, xx, yy;
```

Domain *d;

{

Thread *t;

cell_t c;

int c_max,c_min;

FILE *f1;

int i;

real sum=0.;

```
f1=fopen("interface.txt","a");
```

min=1000.0;

```
max=-1000.0;
```

```
d = Get_Domain(1);
```

thread_loop_c(t,d)

{

```
begin_c_loop(c,t)
```

{

```
if((c<123)&&(c>3))
```

{

 $C_UDSI(c,t,1)=C_UDSI_G(c,t,0)[0];$

```
C_UDSI(c,t,2)=C_UDSI_G(c,t,1)[0];
```

}

```
if ((C_UDSI(c,t,2)>max)&&(C_UDSI(c,t,0)<xb0*1.1)) max=C_UDSI(c,t,2);
```

```
if ((C_UDSI(c,t,2)<min)&&(C_UDSI(c,t,0)>xa0*0.95)) min=C_UDSI(c,t,2);
```

}

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Biographical Information

Md Noor E Alam Siddique was born in Khulna, Bangladesh and got his primary education there. He received his Bachelor of Science (BS) degree in Materials and Metallurgical Engineering from Bangladesh University of Engineering and Technology, Dhaka, Bangladesh in 2007. As recognition for being ranked first in his class in BS program, he was appointed as lecturer in the Materials and Metallurgical Engineering department (MME) to teach undergraduate courses. Also he received Amir Khushbahar Gold medal for holding highest CGPA (grade point) in the MME dept, and Engineers 6061 gold medal for achieving highest CGPA in the Engineering faculty consisting of some other engineering departments. He also received his Master of Science (MS) degree from the same institution in 2009.

In 2009 he joined The University of Texas at Arlington to pursue his doctoral degree and worked under supervision of Dr. Fuqiang Liu of Materials Science and Engineering and received his Ph.D. degree in 2014. His primary research area during his doctoral study was science, engineering and technology of lithium ion battery materials, while he also worked on PEM fuel cell catalyst, vanadium redox flow battery and synthesizing nanostructures for lithium ion batteries. His research results were published in several journals and are listed beleow.

 N A Siddique, Ashley M Allen, Partha P Mukherjee, Fuqiang Li, "Simulation of effect of interfacial lithium flux on miscibility gap in non-equilibrium phase transformation of LiFePO₄ particles", Journal of Power Sources 245 (2014) 83-88. N. Siddique, Amir Salehi, Fuqiang Liu, "Stochastic reconstruction and electrical transport studies of porous cathode of Li-ion batteries", Journal of Power Sources 217 (2012) 437-443.

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He plans to work in the industries to gain more practical experience that he can use later on in his research either in some universities or research institutes. He also plans to start business in future in the developing countries. But mostly he is keen to remain attached to academia somehow.